

JGR Solid Earth

RESEARCH ARTICLE

10.1029/2018JB016773

Special Section:

Merging geophysical, petro-chronologic and modeling perspectives to understand large silicic magma systems

Key Points:

- A new high-resolution numerical thermal model is presented, operating over broad spatial (≥ 1 m to ≥ 10 km) and temporal scales (<100 to $>10^6$ years)
- The time scale over which H_2O -rich fluid exsolves from a crystallizing basaltic sill matches that for transient heating of adjacent wall rock
- Once ambient crust is $\sim 600^\circ\text{C}$, large melt fractions in wall rock (granitoid and mafic sills) adjacent to new basaltic sills may form and advect

Supporting Information:

- Supporting Information S1

Correspondence to:

M. A. Calogero,
calogero@umich.edu

Citation:

Calogero, M. A., Hetland, E. A., & Lange, R. A. (2020). High-resolution numerical modeling of heat and volatile transfer from basalt to wall rock: Application to the crustal column beneath Long Valley caldera, CA. *Journal of Geophysical Research: Solid Earth*, 125, e2018JB016773. <https://doi.org/10.1029/2018JB016773>

Received 10 OCT 2018

Accepted 20 FEB 2020

Accepted article online 21 FEB 2020

High-Resolution Numerical Modeling of Heat and Volatile Transfer from Basalt to Wall Rock: Application to the Crustal Column beneath Long Valley Caldera, CA

M. A. Calogero¹ , E. A. Hetland¹, and R. A. Lange¹

¹Department of Earth and Environmental Sciences, University of Michigan, Ann Arbor, MI, USA

Abstract We present a high-resolution numerical model of the thermal evolution of the crustal column beneath Long Valley caldera, California, from which $>800\text{ km}^3$ rhyolite erupted over the last 2.2 Myr. We examine how randomly emplaced basaltic sills of variable thickness (10, 50, or 100 m) at various depth intervals (10–25 km) and at variable emplacement rates (5–50 m/kyr) gradually heat the crust and lead to a variably mixed crustal lithology (solidified mafic sills and preexisting granitoid). We additionally explore the time scales over which dissolved water ($\sim 3\text{ wt\%}$) in a newly emplaced basaltic sill exsolves during crystallization and is transferred to adjacent wall rock that is undergoing partial melting. We employ a finite-difference-based technique, with variable spatial (≥ 1 m to ≥ 10 km) and temporal (<100 and $>10^6$ years) resolution, that enables dense analysis within and directly adjacent to a newly emplaced sill. Our results show that once ambient crustal temperatures reach ~ 500 – 600°C , subsequent injections of basaltic sills lead to significant partial melting of adjacent wall rock (granitoid and solidified mafic sills) on time scales (10^1 – 10^2 years) that match those of exsolution of H_2O -rich fluid from basaltic sills. Large volumes of fusion ($>10\%$) during fluid-undersaturated partial melting, combined with the preexisting occurrence of aplite dikes, facilitates the development of melt-filled fractures that exceed the critical length for self-propagation. The advection of wall-rock partial melts (with a combined mantle-derived and crustal geochemical signature) to shallower depths will alter both the thermal and compositional profile of the middle-upper crust.

1. Introduction

Understanding the origin and evolution of voluminous rhyolitic magmas is of considerable interest because their formation fundamentally reconstitutes and differentiates continental crust, and they may lead to future “supervolcano” eruptions. The largest explosive eruptions (more than hundreds of cubic kilometers) that have occurred in the United States over the last 1 Myr were high- SiO_2 (75–77 wt%) rhyolites that erupted from the Long Valley caldera in eastern California (e.g., Hildreth & Wilson, 2007) and the Yellowstone caldera in Wyoming (e.g., Christiansen, 2001), both of which are located in the Basin and Range extensional province of the United States (Figure S1 in the supporting information).

Several previous studies have addressed the origin of silicic melts by modeling the thermal evolution of the crust as a function of basalt influx (e.g., Shaw, 1980; Huppert & Sparks, 1988; Petford & Gallagher, 2001; Jackson et al., 2003; Dufek & Bergantz, 2005; Annen et al., 2006; Leeman et al., 2008; Karakas & Dufek, 2015; Colón et al., 2018). These studies examined how different rates of basalt emplacement heat the crust, leading to the onset of crustal melting, and over what spatial and temporal scales this process might occur. These models additionally track the cooling and crystallization of basaltic intrusions, and thus the geochemical evolution of differentiated mantle melts. Some studies have focused on the emplacement of basalts in the lowermost crust (e.g., Annen et al., 2006; Dufek & Bergantz, 2005; Jackson et al., 2003), which is relevant to subduction zone settings, whereas others have focused on extensional environments where basalts ascend higher into the middle crust (e.g., Colón et al., 2018; Karakas & Dufek, 2015; Leeman et al., 2008). There is general consensus that both interstitial melts from the basalt intrusions and partial melts of preexisting crust contribute to the formation of silicic melts but in a complex, multistage process. For example, in the case of the Yellowstone rhyolites, Simakin and Bindeman (2012) and Bindeman and Simakin (2014) have modeled the late-stage, shallow (<10 km) remelting of

hydrothermally altered crust that occurred to form many of the low- $\delta^{18}\text{O}$ Yellowstone rhyolites. Conversely, Leeman et al. (2008) addressed the initial emplacement of basalt into the middle crust (10–30 km) that produced evolved Yellowstone melts from both preexisting crust and mantle-derived basalt. More recently, Colón et al. (2018) examined the crustal-scale consequences of basalt emplacement/entrapment in the middle crust beneath Yellowstone.

The purpose of this study is to examine how the process of basalt emplacement into the granitoid middle to upper crust beneath the Long Valley caldera, California, drives the production of hydrous rhyolitic melts. We build on the crustal-scale approach taken in previous studies (e.g., Annen et al., 2006; Colón et al., 2018; Dufek & Bergantz, 2005; Karakas & Dufek, 2015; Leeman et al., 2008), where the thermal evolution of the entire crustal column in response to basalt influx is simulated. However, this study differs from previous work by examining key processes that occur at relatively short length and time scales (on the order of meters and years, respectively), which requires a numerical model that can achieve such high spatial and temporal resolution while still being able to simulate the entire lithospheric column over a time scale of millions of years.

Our goal in this study is to pay particular attention to the formation of localized partial melts along the wall-rock margins of basaltic sills within the first tens to hundreds of years of their emplacement (e.g., Lim et al., 2012), even when long-term (>1 kyr) ambient temperatures remain sub-solidus (i.e., $<650^\circ\text{C}$). This is a process that has not yet been quantitatively modeled in the literature and may contribute to the formation, segregation, and ascent of evolved partial melts of a mixed lithology (granitoid and mafic sills), leading to their gradual accumulation in the upper crust. These partial melts will have a combined mantle-derived and crustal geochemical signature. The time scale for the segregation and ascent of these localized partial melts is also considered.

Another focus of this study is the examination of how water is transferred from basaltic sills to partial melts of the crust. One of the defining features of the voluminous ($>400\text{ km}^3$) Early Bishop rhyolite is its high concentration of water ($\leq 6.5\text{ wt\%}$; Wallace et al., 1999; Anderson et al., 2000). The source of this water cannot be the preexisting granitoid crust, with a bulk water content $<0.6\text{ wt\%}$ held in biotite and hornblende (Bateman, 1992), and instead must be derived from the Long Valley basalts (e.g., Cousens, 1996), which erupted in close spatial and temporal association with the rhyolites of Long Valley caldera. Understanding this transfer of water is of major importance because melt fraction-temperature (F-T) curves depend strongly on how much water is present.

2. Tectonic/Geologic Setting for Numerical Modeling

The numerical simulations presented in this study are designed to be applicable to the Long Valley magmatic system, located in eastern California along the western edge of the Basin and Range province (Figure S1b). Between 2.2 and 0.8 Ma, $\sim 100\text{ km}^3$ of rhyolite erupted to form the Glass Mountain complex (Hildreth, 2004; Metz & Mahood, 1991), which was followed by the climactic eruption of the $>600\text{ km}^3$ Bishop Tuff within ~ 6 days (Wilson & Hildreth, 1997) to form the Long Valley caldera at $\sim 765\text{ ka}$ (Andersen et al., 2017). Soon after, the eruption of $\sim 100\text{ km}^3$ of rhyolite (named “Early” post-caldera rhyolite by Bailey, 1989) occurred within the caldera between ~ 750 and 640 ka (Hildreth et al., 2017). Over the last 600 kyr, $\leq 8\text{ km}^3$ of rhyolite have erupted inside the margins of the Long Valley caldera (Hildreth, 2004; Simon et al., 2014). A combined volume of $>800\text{ km}^3$ of rhyolite has erupted within the Long Valley volcanic system over the last 2.2 Myr but over discrete, punctuated time periods characterized by highly variable eruption rates.

This voluminous production of rhyolite was driven by an influx of basalt into the crustal column beneath Long Valley caldera (e.g., Hildreth & Wilson, 2007). Seismic studies indicate a mid-crustal (~ 10 - to 20-km depth) low-velocity zone beneath Long Valley (Flinders et al., 2018; Frassetto et al., 2011; Thurber et al., 2009), consistent with the presence of melt and/or fluid. The Long Valley basalts were formed by partial melting of subduction-modified mantle lithosphere (Cousens, 1996), which was triggered by upwelling asthenosphere due to lithosphere delamination beneath the eastern Sierra Nevada (e.g., Manley et al., 2000; Saleeby et al., 2003; Zandt et al., 2004).

3. Numerical Model

Temperature and melt fraction are solved for using the 1-D nonhomogeneous transient thermal diffusion equation:

$$\rho C_p \frac{\partial T}{\partial t} + \rho L \frac{\partial F}{\partial t} = \frac{\partial}{\partial x} k(x) \frac{\partial T}{\partial x}, \quad (1)$$

where ρ is density, C_p is specific heat capacity, L is the latent heat, $k(x)$ is thermal conductivity, T is temperature, F is melt fraction, t is time, and x is depth. Temperature and melt fraction are related by a piecewise linear function, $F(T)$, given in general as

$$F(T) = \begin{cases} 0, & T \leq T_0, \\ m_i T + b_i, & T_{i-1} \leq T < T_i, \\ 1, & T \geq T_p, \end{cases} \quad (2)$$

where $T_0 = T_{\text{solidus}}$ and $T_p = T_{\text{liquidus}}$, and m_i and b_i are the slope and intercept, respectively, for $i \in [1, p]$ segments of the temperature versus melt fraction relationship. Approximating this relationship with $F(T)$ allows for the general equation

$$\frac{\partial F}{\partial t} = \phi \frac{\partial T}{\partial t}, \quad (3)$$

where $\phi = 0$ if $T < T_{\text{solidus}}$ or $T > T_{\text{liquidus}}$, and $\phi = m_i$ otherwise. Using equation (3), equation (1) then reduces to

$$(\rho C_p + \phi \rho L) \frac{\partial T}{\partial t} = \frac{\partial}{\partial x} k(x) \frac{\partial T}{\partial x}. \quad (4)$$

The 1-D compaction and expansion due to solidification and melting, respectively, are approximated by

$$\frac{dX}{dt} = \theta X \frac{dF}{dt}, \quad (5)$$

where X is the height of the column of the region undergoing solidification or melting, and θ is a scalar representing the percent height change during the phase transition.

Equation (4) is solved in nondimensional form using a numerical method of lines (MOL) technique on a nonuniform spatial grid and with adaptive time resolution. Discretizing the spatial coordinate as x_j , $j \in [1, N]$, we denote $T(x_j, t) = T_j$, which for brevity is not written as a function of t . The spatial finite difference approximation of equation (4) is

$$\frac{\partial T}{\partial t} \simeq \frac{1}{\rho(C_p + L\phi_i)} \frac{\left(\frac{T_{j+1} - T_j}{x_{j+1} - x_j}\right)k_{j+\frac{1}{2}} - \left(\frac{T_j - T_{j-1}}{x_j - x_{j-1}}\right)k_{j-\frac{1}{2}}}{\frac{1}{2}(x_{j+1} - x_{j-1})}, \quad (6)$$

where $k_{j+\frac{1}{2}} = \frac{k_{j+1} + k_j}{2}$ and $k_{j-\frac{1}{2}} = \frac{k_{j-1} + k_j}{2}$. Equation (6) is numerically integrated at x_j using a standard Runge-Kutta method with adaptive time steps (e.g., Dormand, 1996). In the adaptive time-stepping method, the nodal temperatures at each time step are calculated using a proposed time step with both second- and third-order Runge-Kutta integration schemes. If the maximum difference in the solutions is larger than a specified tolerance (here we use 10^{-5} in nondimensional temperature), the current time step is decreased, and the temperatures for the next time step are recalculated using the two Runge-Kutta schemes and the smaller time step. We continue to decrease the time steps until the second- and third-order solutions are within tolerance. For each rejected time step, the time step is decreased to $0.9 \text{del}T (\text{tolerance/error})^{1/3}$, where error is the maximum absolute difference in the temperature solutions and $\text{del}T$ is the previously tested time step (e.g., Hetland et al., 2010; Noda et al., 2009). Once the solutions converge within tolerance, we take the third-order solution as the temperatures for the next time step, and the next time step is also $0.9 \text{del}T$.

Table 1
Material Properties

	Symbol	Mantle	Lower crust	Granitoid	Basalt
Density, solid (kg/m ³)	ρ_S	3,400	3,050	2,650	3,100
Density, melt (kg/m ³)	ρ_M			2,650	2,830
Specific heat (J/(kg · K))	C_p	1,131	1,390	1,370	1,480
Latent heat (J/kg)	L	0	3.5E+05	2.7E+05	4.0E+05
Conductivity (J/[s · m · K])	k	3.4	2.6	3	2.6
Melt compaction		0	0	0	0.1

Note. References: ρ , Holbrook et al. (1992) and Kay et al. (1992); C_p and L , Bohrsen and Spera (2001); k , Chapman and Furlong (1992) and Hofmeister (1999).

(*tolerance/error*)^{1/3} (i.e., if the error is larger than 1.372 of the tolerance, the time step grows). The minimum nodal spacing controls the stability of each time step, so the solution is always within stability for the larger nodal spacings on the nonuniform spatial grid. Using this scheme, the time step decreases when the temperature gradients are high and then gradually increases after the spatial variations in the temperature gradients subsequently subside. Once the temperatures are calculated at a given time step, the nodal melt fraction is solved for using equation (2). The melt fraction is then used to calculate the material density, using prescribed densities of the solid and the liquid phases. A first-order discretized form of equation (5) is integrated to find the compaction (and/or expansion) due to changes in melt crystallization (and/or melt formation).

In all of the simulations presented in this study, the initial domain thickness (from the surface to depth) is 45 km. Over this domain, N nodes are initially distributed evenly with depth. The material properties of each node are assigned based on their depth, such that if a node is within the boundaries of the lower crust, it is assigned material properties of the lower crust, etc. (Table 1). The first sill is then injected into the model domain (all basalt sills in this study have an initial temperature of 1150 °C and melt fraction of 1), and then the nodes are adjusted as described below, before solving for the temperatures. Temperatures are calculated at each time step as described above, with additional basalt sills periodically injected into the crust at specified emplacement rates.

After each sill is injected, the domain thickness is increased by the thickness of the sill, and the temperature and melt fraction fields below the sill are advected downward. With each emplacement, new nodal positions are calculated using a node placement optimization motivated by a moving mesh MOL (MM-MOL) technique (Huang & Russell, 2011). This allows temperatures to be resolved where thermal gradients are high, without overly discretizing regions of the domain with near-linear temperature gradients. In this node adaptation scheme, the spatial resolution is increased in regions of the model domain with high curvature of the temperature field and/or with high melt fraction, while node spacing is decreased where the temperature has minimal curvature and there is no melt present. To adapt the node distribution, we use the mesh density function

$$\delta^* = 1 + |\nabla^2 T|^{\frac{1}{4}} + \gamma F \quad (7)$$

and solve for optimal nodal spacings that are proportional to δ^* , under the constraint of a prescribed minimum node spacing and that the entire domain is covered by nodes. The proportionality factor γ is determined such that the computation domain contains a fixed number of nodes (i.e., nodes are not added or removed from the domain). To avoid singularities in the temperature field, we smooth the mesh density function, following the MM-MOL method (Huang & Russell, 2011). Compression (expansion) during crystallization (melting) is then accounted for by perturbing the nodal positions. Finally, the node positions may be slightly perturbed in order to ensure that there is at least one node between all material boundaries, including locations of all past sills. After the node positions are adjusted, the temperature and melt fraction fields, as well as all material parameters, are mapped from the old mesh to the new mesh using a linear interpolation.

The numerical solution is implemented in Matlab. With 600 total nodes over an initial domain thickness of 45 km, the largest node spacing after a sill injection is about 130 m and occurs at the edges of the model

domain (close to the surface and the base of the domain). Aside from some resolution tests, the main simulations presented in this study use a minimum spatial resolution of 1 m, which corresponds to a minimum temporal resolution of about a minute. Using a spatial resolution of at least 5 m, the smallest time step is ~6 days. Due to the adaptive time-stepping algorithm, as temperatures diffuse and basalt sills crystallize, temporal resolution increases, and 1,000 years after a sill injection, the time step increases to ~10 days and ~8 months, respectively, for 1- and 5-m minimum spatial resolutions. Simulations of the thermal evolution of the crust after 200 sill injections take ~14 and ~2 hr for 1- and 5-m minimum spatial resolutions, respectively, on a circa 2012 desktop workstation. Results from a benchmark simulation of basalt intruding basalt and subsequent cooling is provided in Figure S2.

3.1. Model Parameters

An initial 20 °C/km linear geotherm is assumed prior to injection of the basaltic sills, which is consistent with the seismically imaged lithosphere-asthenosphere boundary at ~60 km for the Basin and Range (Levander & Miller, 2012), where a temperature of 1200 °C marks the lithosphere-asthenosphere boundary (LAB). This relatively cool geotherm of 20 °C/km is also consistent with the volcanic hiatus in the Long Valley region from 8 to 4.5 Ma prior to renewed magmatism in the region (Bailey, 1989; Du Bray et al., 2016). The initial crustal composition is assumed to be granitoid from the surface to 32-km depth, mafic amphibolite from 32–35 km depth, and mantle lithosphere below 35 km (Figure S3; Fliedner et al., 2000). The bulk composition used for the granitoid upper crust beneath Long Valley caldera is the average composition of upper crust from Rudnick and Gao (2014). Note that several whole-rock compositions of granitoid from the Sierra Nevada batholith (e.g., Bateman, 1992; Coleman et al., 2012) broadly overlap this average upper crust composition. The composition of basalt used in this study is the average of several Long Valley basalts (50.9 wt% SiO₂, 1.5 wt% TiO₂, 15.7 wt% Al₂O₃; 8.2 wt% FeO_T, 8.0 wt% MgO, 9.1 wt% CaO, 3.6 wt% Na₂O, 2.2 wt% K₂O, and 0.6 wt% P₂O₅).

In this study, we initially focus on the emplacement of 5 km of basalt over a 1-Myr time period, in the form of 50-m thick sills emplaced every 10 kyr, into the crust (5 m/kyr). This rate of emplacement was also employed in Leeman et al. (2008) in their thermal model of the crust in response to Yellowstone hot spot magmatism. Different sill thicknesses (10 and 100 m) and different time intervals for the emplacement of basalt are additionally explored, leading to different basalt flux rates (~5–50 m/kyr). Because rhyolite eruption rates out of the Long Valley volcanic system have been highly variable, from ~0.1 km³/kyr between 2.2 and 0.8 Ma (Hildreth, 2004) to >600 km³/week (Wilson & Hildreth, 1997), the effect of variable basalt emplacement rates, including relatively high rates over short time periods (e.g., 50 m/kyr for ≤100-kyr duration), merits evaluation.

We prescribe a sill emplacement history, either placing sills at the same fixed depth or randomly placing them within a depth interval of the upper crust. Because the tectonic setting is continental extension (and not subduction), basaltic sills are emplaced in the middle (not lower) crust. Note that mid-crustal depth intervals (10–25 km) are broadly consistent with seismic evidence of a mid-crustal sill complex beneath the Snake River Plain due to the influx of basalt from the Yellowstone plume (Stachnik et al., 2008; Yuan et al., 2010), which in turn drove the formation of voluminous rhyolite along the Yellowstone hot spot track (Colón et al., 2018; Leeman et al., 2008). In all of the simulations presented in this paper, we assume that all injected basalt sills have an initial uniform temperature of 1150 °C and are emplaced at fixed time intervals. Nodes within the sills, both the most current and previously solidified sills, are assigned the material properties of basalt (Table 1), and these properties, apart from density, do not change after the sills have solidified. Nodes outside of the sills are assigned material properties of the wall rock (Table 1), relative to the evolving material boundaries with each emplacement.

When sills are emplaced, all material boundaries below the sill margin are moved down by the sill thickness. In the case where sills are inserted at a fixed depth, the top of each new sill is at the same depth, and previously emplaced sills are pushed down, creating a contiguous depth range of uniform basaltic composition. When sills are inserted randomly within a depth range, a mixed lithology of alternating granitoid and solidified mafic sills is created. The minimum nodal spacing of 1 m ensures that a new sill of 50 m is resolved with 51 nodes as it solidifies. Hereafter, in this text, the term “basaltic sill” refers to a newly emplaced sill that is undergoing crystallization, whereas the term “mafic sill” refers to a previously solidified basaltic sill.

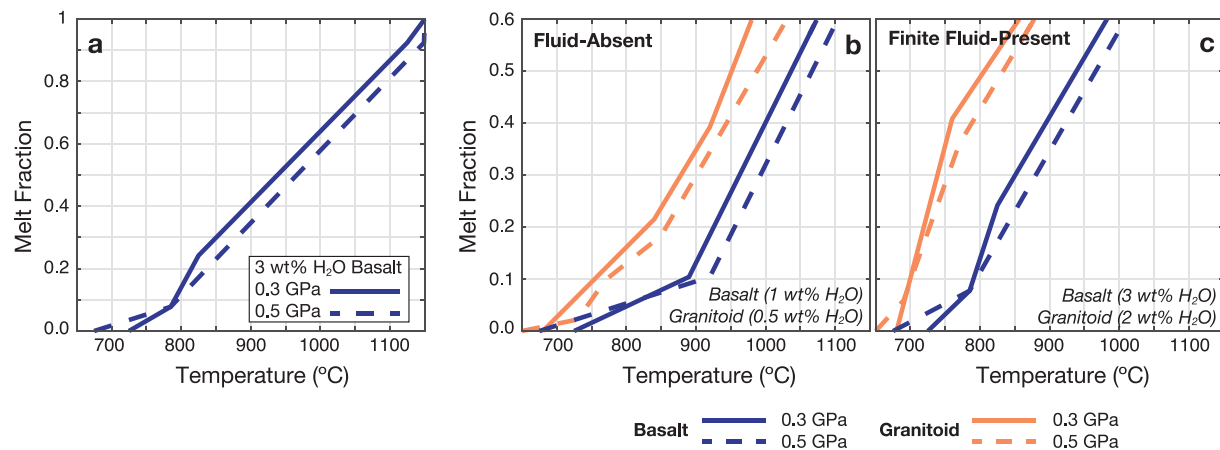


Figure 1. Melt fraction versus temperature (F-T) relationships for mafic sills and granitoid, calculated with rhyolite MELTS (Gualda et al., 2012). (a) Basalt crystallization at 0.3 GPa (solid) and 0.5 GPa (dashed). (b) Case 1: Fluid-absent granitoid (orange) and mafic sill (blue) melting. (c) Case 2: Finite fluid-present granitoid and mafic sill melting.

3.2. Basalt Crystallization F-T Curve

We assume an initial temperature of 1150 °C and an initial water concentration of 3 wt% H₂O for each basaltic sill. This water content is based on analyses of ≤ 3 wt% H₂O in olivine-hosted melt inclusions of nearby Pliocene–Quaternary basalts erupted along the eastern margin of the Sierra Nevada range (Gazel et al., 2012). To determine the basalt equilibrium crystallization curve with temperature, the thermodynamic model rhyolite-MELTS (Gualda et al., 2012) was used at 0.3 and 0.5 GPa. The results are shown in Figure 1a and are parameterized into four linear segments. According to the MELTS calculation, latent heat is distributed nearly linearly as a function of degree of crystallization for this melt composition, which allows total enthalpy (L) to be treated as a linear function of melt fraction. The effect of a non-linear release of latent heat across the liquidus-solidus interval is explored in section 4.1.3.

3.3. F-T Curves for Crustal Melting

Melt fraction versus temperature curves are needed for both solidified mafic sill and granitoid lithologies, and they will each depend strongly on the total amount of H₂O present. Two end-member H₂O conditions for each lithology are evaluated: fluid-absent and finite fluid-present (i.e., fixed amount of H₂O present; not unlimited). In all cases, the F-T curves are calculated using the rhyolite-MELTS thermodynamic model of Gualda et al. (2012) and are plotted in Figure 1.

3.3.1. Case 1: Fluid-Absent Mafic Sills and Granodiorite

In the first case, it is assumed that after complete crystallization of the basaltic sills in the granitoid crust, approximately one third (1 wt% H₂O) of the initial 3 wt% H₂O dissolved in the melt is retained in the mineral assemblage in the form of biotite and/or hornblende, which leaves ~ 2 wt% H₂O as a fluid phase under sub-solidus conditions. In this first-case scenario, it is assumed that this fluid phase is completely lost from the sill and surrounding crustal rocks (i.e., through ascent along grain boundaries and/or fractures). This leaves a fluid-absent mafic sill with a bulk H₂O content of ~ 1 wt%. Similarly, it is assumed that the surrounding granitoid rock (with a bulk composition similar to that of average upper crust; Table 1) is also devoid of any fluid phase, and H₂O is only present in the form of biotite and hornblende ($< 10\%$ of each phase; Bateman, 1992). In this case, a total H₂O content of ~ 0.5 wt% is assumed for the granodiorite wall rock. The F-T curves for the fluid-absent mafic sill and granodiorite, with 1 and 0.5 wt% H₂O, respectively, are calculated using rhyolite-MELTS and are illustrated in Figure 1b. Note that despite the lower H₂O content, the fluid-absent granodiorite has a higher melt fraction relative to the fluid-absent mafic sill at any given temperature above the solidus. For example, at 800 °C and 0.3 GPa, the granodiorite and mafic sill are predicted to contain $\sim 15\%$ and $\sim 5\%$ melt, respectively.

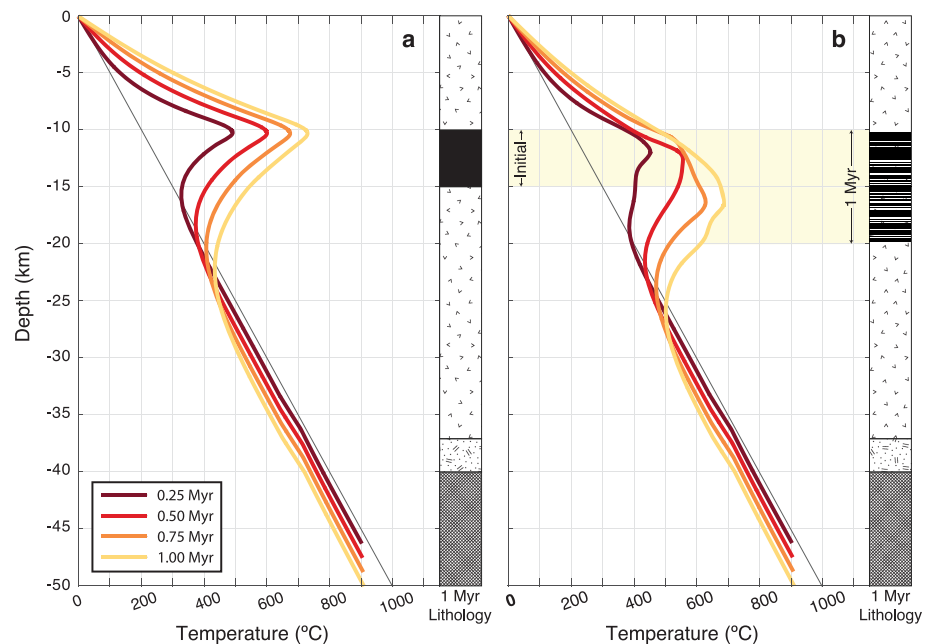


Figure 2. Thermal profiles for (a) fixed-depth emplacement at 10-km depth and (b) random-depth emplacement with an initial depth range of 10–15 km. Emplacement rate is 5 km/Myr in both cases (50-m thick sill every 10 kyr). Profiles are shown after 0.25 Myr (maroon), 0.5 Myr (red), 0.75 Myr (orange), and 1 Myr (yellow). The crustal lithology after 1 Myr for each emplacement scheme is included, where black is basalt sills and the other rock types are defined in the legend in Figure S3. The thin black line denotes the initial geothermal gradient. Fluid-absent F-T curves at 0.3 GPa are used for the wall rock (Figure 1b).

3.3.2. Case 2: Finite Fluid-Present Mafic Sills and Granodiorite

In the second case, it is assumed that during late-stage crystallization of the basaltic sills, ~2 wt% H₂O exsolves as a fluid phase and reaches the interface between the sill margin and the granodiorite wall rock. During transient heating of wall rock adjacent to a newly emplaced sill, an initially fluid-present melting

curve at the solidus for granodiorite is employed, based on a whole-rock composition of 2 wt% H₂O, which becomes fluid undersaturated as melt fraction increases. Also shown is an initially fluid-present melting curve for the solidified mafic sills at the solidus, based on a whole-rock composition of 3 wt% H₂O (2 wt% H₂O fluid + 1 wt% in hydrous mineral phases), which becomes fluid undersaturated as melt fraction increases. A comparison of the two initially fluid-present F-T curves (Figure 1c), which become fluid undersaturated with melt fraction, shows that the granodiorite wall rock produces a far greater melt fraction than the solidified mafic sills at a given temperature (e.g., 30% vs. 5% at 750 °C and 0.5 GPa). We use the term “finite fluid-present” to distinguish the F-T curves in Figure 1c (Case 2), which are based on a finite (i.e., fixed) amount of water in the system, from F-T curves that are fluid saturated under all melt fraction conditions (i.e., unlimited amount of water available).

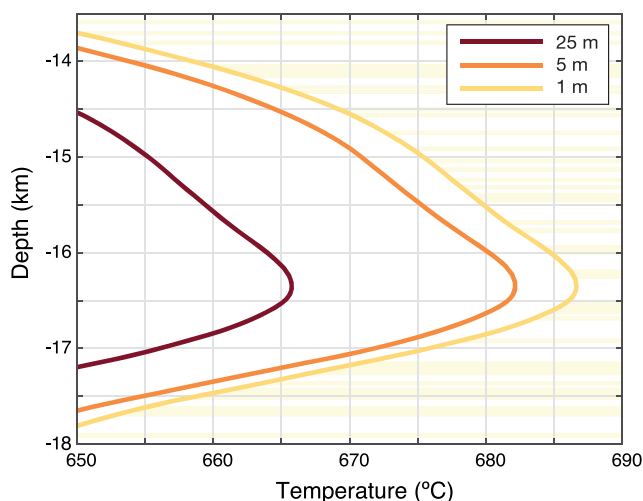


Figure 3. Thermal profiles after 1 Myr of random-depth emplacement (50-m sill every 10 kyr; fluid-absent at 0.3 GPa, Figure 3b) for three minimum spatial resolutions: 25 m (maroon), 5 m (orange), and 1 m (yellow). Pale horizontal background blocks mark the sill positions at 1 Myr. The most accurate results are for the highest spatial resolution (1 m).

4. Results

4.1. Crustal-Scale Thermal Evolution

We ran simulations with the following parameters varied: (1) sill emplacement style (e.g., fixed vs. random emplacement depth); (2) minimum spatial resolution (25, 5, and 1 m); (3) nonlinear

distribution of latent heat in basaltic sills; (4) thickness of sills (10, 50, and 100 m); (5) depth range of sill emplacement; (6) rate of sill emplacement (i.e., different time intervals over which a fixed amount of basalt is emplaced in the crust); and (7) fluid-absent versus finite fluid-present F-T curves for wall rock (granodiorite and solidified mafic sills). In each of these cases, all other variables were held constant. The goal was to evaluate first-order effects on the thermal evolution of the crust at the kilometer length scale, as well as the resulting proportions of mafic sills and granitoid in the crustal column. Both crustal temperature and the extent of a mixed lithology in the crustal column must first be determined before the extent of any partial melting along the margins of newly emplaced sill can be evaluated.

4.1.1. Fixed Versus Random Depth Emplacement

Previous studies have examined both fixed and random emplacement of basaltic sills (e.g., Annen et al., 2006; Leeman et al., 2008). For comparative purposes, we also consider both emplacement schemes. In each simulation, the initial sill temperature (1150 °C) and emplacement rate (50-m sill every 10 kyr) are held constant. In the fixed-depth emplacement scheme, new sills are injected with their top surface at 10-km depth, the shallowest depth chosen in Leeman et al. (2008). In the random-depth emplacement scheme, the initial depth range for sill emplacement is 10–15 km, with the upper depth limit held constant at 10 km and the lower depth limit moved down by the thickness of each new sill after it is emplaced. The simulations were run for 1 Myr in both cases, resulting in 5 km of basalt added to the crustal column. In the fixed-depth emplacement scheme, the added basalt sills are contiguous, whereas in the random-depth emplacement scheme, a mixed lithology results between 10- and 20-km depth, with an average 1:1 ratio of mafic sills to granitoid. The evolution of the thermal profile at 0.25, 0.5, 0.75, and 1 Myr for each emplacement scheme is shown in Figure 2. The thermal profiles shown are those 10 kyr after the most recent sill emplacement (i.e., immediately before the next scheduled emplacement of a 50-m sill). For both emplacement schemes, all basaltic sills are fully crystallized after 1 Myr, and there are no crustal depths hotter than 725 °C.

In the fixed-depth emplacement simulation, after 1 Myr, the maximum crustal temperature is ~725 °C and occurs inside the most recently emplaced basalt sill, immediately below 10-km depth (i.e., near the interface with granitoid crust). In this simulation, temperatures are >600 °C over ~2-km thickness of granitoid crust immediately above the contiguous mass of basaltic sills. In contrast, in the random-depth emplacement simulation, after 1 Myr, the maximum crustal temperature is ~686 °C, located ~0.27 km shallower than the most recently emplaced basalt sill (center at ~16.62 km). The amount of granitoid (distinct from mafic sills) heated >600 °C is nearly twice as thick, at ~4 km. The increase in thickness of granitoid that is heated by the random emplacement of sills reflects the more efficient transfer of heat from the basaltic sills to the surrounding granitoid when the 50-m thick sills are spread out throughout the crustal column (e.g., Dufek & Bergantz, 2005). A microscale example of this phenomenon is described in Spera and Bohrsen (2018).

4.1.2. Variation of Minimum Spatial Resolution

To evaluate the effect of the spatial resolution of our finite-difference model on crustal thermal profiles, we ran the same random-emplacement simulation (i.e., with identical sill history) from Figure 2b, using minimum spatial resolutions of 1, 5, and 25 m. The final temperature profiles over the depth range where most sills were emplaced are shown in Figure 3. For reference, the center of the last sill emplaced is at a depth of 16.62 km. After 1 Myr, the maximum temperatures in the 25- and 5-m cases are ~19 and ~4 °C, respectively, colder than the 1-m case. The most precise results are obtained with the highest spatial resolution (1 m), which is used for all simulations presented in this study.

4.1.3. Varying Distribution of Latent Heat Across the Liquidus-Solidus Interval

In equation (1), if $F > 0$, then the latent heat (L) is released linearly as a function of percent crystallization. This is the default case for all simulations presented in the main body of the text. To test whether crustal-scale thermal profiles are sensitive to how latent heat is distributed across the liquidus-solidus interval of basaltic sills, we added an additional step during the Runge-Kutta time-stepping, which adjusts the proportion of latent heat released across the liquidus-solidus interval (Figure S4a). In the first test case, 75% of the latent heat is released in the first 25% of crystallization, whereas in the second test case, only 25% of the latent heat is released in the first 75% of crystallization, and the remaining 75% of latent heat is released in the last 25% of crystallization. In all simulations, the same total amount of latent heat is released during crystallization of the basaltic sills, and all other parameters are unvaried, including the random sill emplacement history (same as that in Figure 2b). The resulting thermal profiles at 1 Myr are within ≤ 5 °C

for all three cases (Figure S4b), which shows that variations in how latent heat is distributed across the liquidus-solidus interval have a negligible effect on crustal-scale thermal profiles.

4.1.4. Variation in Thickness of Basaltic Sills

The effect of varying the thickness of basaltic sills (10, 50, and 100 m), which are emplaced at an identical long-term rate (5 km/1 Myr) and initial depth interval (10–15 km), on the crustal-scale thermal profile was also explored. There is no difference in the temperature profile through the crust when 10- and 50-m thick sills are randomly emplaced (Figure S5). However, three simulations of random emplacement of 100-m thick sills every 20 kyr were run, producing variable thermal profiles after 1 Myr. For one simulation, the thermal profile closely matches those for the 10- and 50-m thick sill emplacement runs, whereas other random sill histories produced higher temperature concentrations at shallower or deeper depths (see Figure S5b). The differences in the three 100-m sill thermal profiles is due to the random (and thus varying) emplacements of relatively thicker and fewer basaltic sills into the crustal column, particularly controlled by the final few sill emplacements in each simulation.

4.1.5. Variation of Emplacement Depths and Depth Ranges

In order to test the effect of varying the depth ranges over which basaltic sills are emplaced, we ran simulations of random sill emplacements (50 m/10 kyr) over a 1-Myr time interval with variable initial depth ranges: (1) 10–15 km, (2) 15–20 km, and (3) 10–20 km. In each of the three test cases, we ran four realizations of random sill emplacements and took their average thermal profile after 1 Myr (Figure S6). In all cases, a total of 5 km of basalt was emplaced in the respective depth ranges after 1 Myr (final ranges: 10–20 km, 15–25 km, and 10–25 km, respectively). In the models where the initial depth range spans 5 km (10–15 and 15–20 km), the final sill:granitoid ratio is 1:1, whereas when the initial depth range spans 10 km (10–20 km), the resulting sill:granitoid ratio is 1:2. The most important observation in Figure S6 is that the depth interval over which basaltic sills are emplaced strongly affects the proportion of mafic sills embedded within the granitoid crust once temperatures reach $\sim 600^\circ\text{C}$. In other words, the deeper and broader the initial depth interval of sill emplacement (at a constant rate), the lower the amount of basalt that is required to bring crustal temperatures to $\sim 600^\circ\text{C}$.

4.1.6. Variation of Sill Emplacement Rate

We next explored the thermal effect of emplacing the same 5-km thickness of basalt over different time intervals (1, 0.5, and 0.1 Myr). In all cases, the initial depth range is 10–15 km, and after 100 sills are emplaced, the average final sill:granitoid ratio is 1:1 between 10- and 20-km depth. The only difference among the three simulations is the frequency with which the 50-m thick sills are emplaced, every 10, 5, and 1 kyr, respectively.

The results for these three simulations are compared in Figure 4, using the same sill injection histories shown in Figure 2b (100 total sills). The lowest basalt emplacement rate (5 km over 1 Myr) results in a thermal profile that is everywhere $<700^\circ\text{C}$, with a maximum temperature of 686°C 10 kyr after the last sill is emplaced. In the second simulation, the same 5 km of basalt is emplaced in half as much time (500 kyr), and the thermal profile is notably hotter, with ~ 4 km of crust reaching temperatures $>700^\circ\text{C}$, 5 kyr after the 100th sill is emplaced, and a maximum at 752°C . In the third simulation, where 5 km of basalt is emplaced within 100 kyr (50 m/1 kyr), ~ 7.7 km of crust is heated above 700°C , 1 kyr after the 100th sill is emplaced, including ~ 2 km in excess of 800°C ; the maximum temperature is 836°C .

These combined results (Figure 4) show that the time interval over which a fixed amount of basalt is placed into the crust has a notably strong effect on the thermal profile through the crust, as found in previous studies (e.g., Annen et al., 2006; Dufek & Bergantz, 2005; Leeman et al., 2008). Over shorter time intervals, there is insufficient time for the heat from the basaltic sills to diffuse away, allowing the crustal range with a 1:1 basalt:granitoid lithology to attain markedly higher temperatures. Conversely, less basalt (3 vs. 5 km) is required at high versus low emplacement rates (50 vs. 5 m/kyr) to reach a similar temperature profile in the crust (Figure 4b, dashed).

4.1.7. F-T Curves for Granitoid Wall Rock: Fluid-Absent Versus Finite Fluid-Present

The final parameter that we explored is the effect of variable melt-fraction versus temperature curves for the wall rock under fluid-absent (Case 1) and finite fluid-present (Case 2) conditions. The latent heat of fusion term in equation (4) during partial melting acts as a heat sink (opposite to the release of heat during basalt crystallization), and hence the choice of the F-T curve for the wall rock can affect crustal-scale thermal profiles. The magnitude of this effect is illustrated in Figure 5 after 100 sill emplacements at two rates, 50 m/10

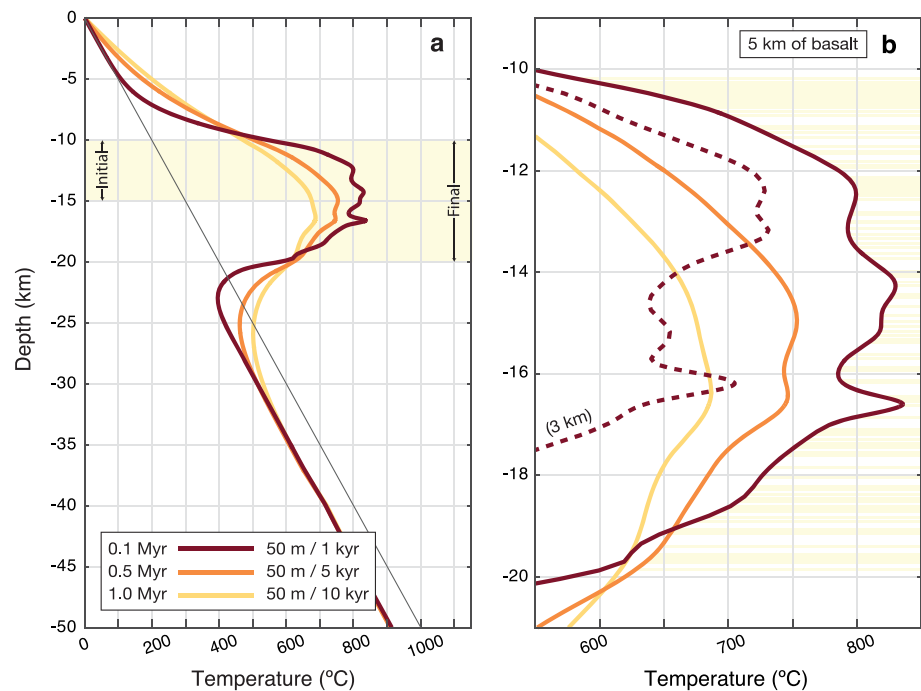


Figure 4. Thermal profiles for different rates of random emplacement: 5 km over 1 Myr (yellow), 0.5 Myr (orange), and 0.1 Myr (maroon). (a) Thermal profile over 0- to 50-km depth with initial and final emplacement ranges marked in yellow. (b) Zoomed-in thermal profile (9–21 km depth). Yellow shading denotes all sill positions after 5 km of basalt have been emplaced. All three simulations use the same sill position history and fluid-absent F-T curves at 0.3 GPa (Figure 3b). At depths above and below the sill emplacement range, the crustal temperatures are colder when emplacement rates are higher, because there is less time for heat from the basalt to diffuse into cooler regions. This is also why the thermal profile is less smooth and more variable over the 10- to 20-km depth interval for the shortest time interval tested (100 kyr).

kyr and 50 m/5 kyr (1 and 0.5 Myr, respectively). For the same amount of basalt injected into the crust (5 km), the lower versus higher emplacement rates heat ~4 versus ~7 km of crust above 650 °C (Figure 5a).

In these simulations, there is a marked difference in the amount of partial melt that is formed, depending on which F-T curve is applied to the wall rock (fluid-absent vs. finite fluid-present), examined in Figure 5b. For the lower emplacement rate, the difference in the melting curves leads to the presence (~2% melt) versus absence of melt. For the higher emplacement rate (i.e., hotter crustal temperatures), the fluid-absent melting curve leads to the formation of less partial melt (<5%) across 3.6 km of crust, whereas the finite fluid-present melting curve leads to the formation of partial melt across 4.5 km of crust, with >20% melt fraction present over nearly 3 km in granitoid wall rock and <2.5% melt fraction present in mafic sills. Importantly, the difference in the amount of partial melt that is formed affects the crustal-scale thermal profile. The latent heat of fusion term causes cooling, the magnitude of which is dependent on the amount of partial melt that is formed. For the high-emplacement rate scenario (Figure 5b), the difference between ~5% and ~25% melt present (fluid-absent vs. finite fluid-present F-T curves) in the crustal column leads to an average cooling of ~15 °C.

4.2. Sill-Scale Thermal Evolution

To evaluate the thermal evolution of the basaltic sills and adjacent wall rocks at high spatial (1–50 m) and temporal (tens to hundreds of years) resolution, simulations were run using the same sill-emplacement history shown in Figure 2b and each of the F-T curve cases introduced in section 3.3 (Figure 1). The objective was to evaluate changes in the temperature and melt fraction of the adjacent wall rock within 50 m of the basaltic sill margin and during the first few hundred years after the sill is emplaced. In all simulations, the emplacement rate is 50 m/10 kyr. The following key factors were varied: (1) ambient temperature of the wall rock immediately prior to each sill emplacement, which increases on average as the simulation progresses (Figure S7), (2) wall rock lithology (granitoid vs. solidified mafic sill), which becomes increasingly variable due to the random emplacement scheme (Figure 6), (3) fluid-absent versus finite

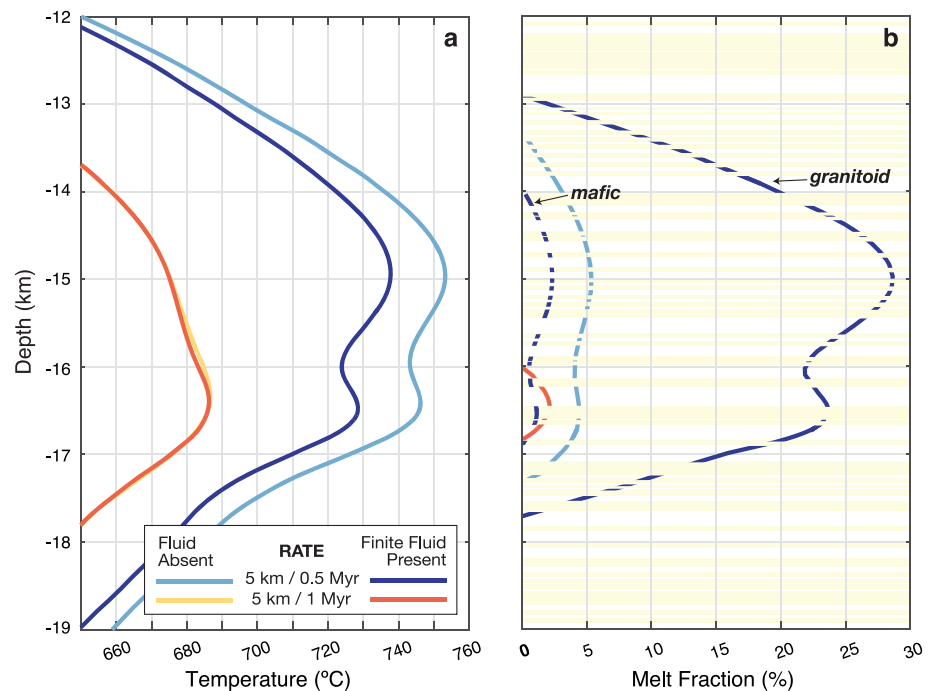


Figure 5. (a) A comparison of thermal profiles for two different emplacement rates: 5 km in 1 Myr (red/yellow) and 0.5 Myr (blues) for fluid-absent (light blue/yellow) and finite fluid-present conditions (dark blue/red). The results show the higher temperatures that result at higher emplacement rates and that fluid-absent versus finite fluid-present F-T curves (0.3 GPa, Figure 3) exert an influence. (b) Melt fraction profiles for the thermal profiles shown in (a). When melt fractions are low, there is little effect on thermal profiles. When melt fractions are high (dark blue vs. light blue), there is marked cooling ($\leq 20^\circ\text{C}$) of the thermal profile due to the effect of latent heat of melting for different melt fractions ($\leq 5\%$ vs. 25%). See text for more details.

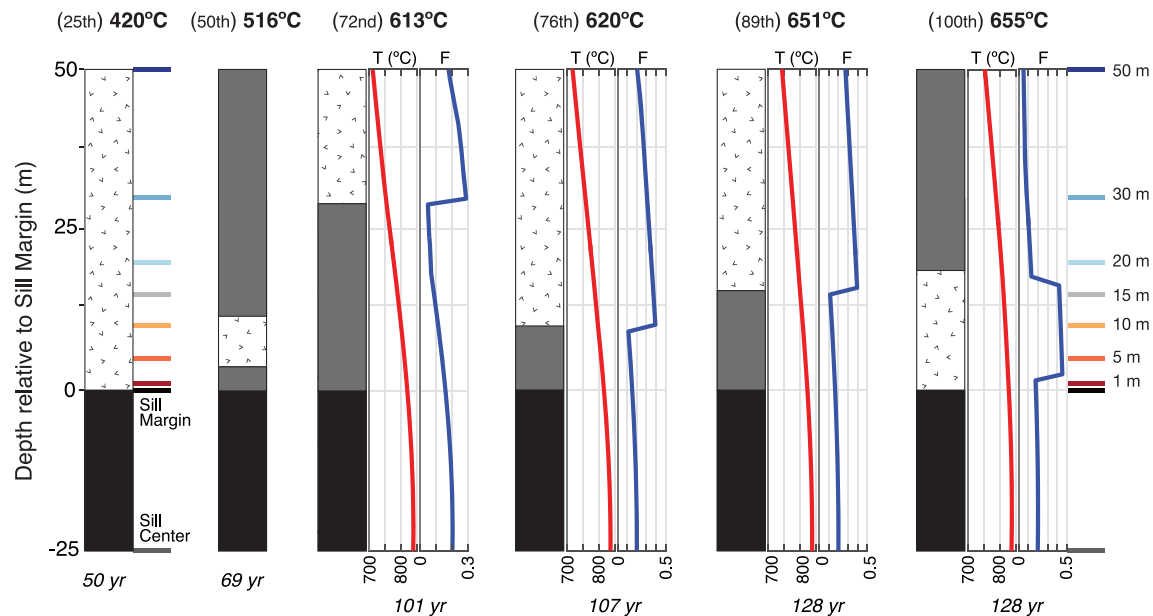


Figure 6. Examples of resultant lithologic column 0–50 m above newly emplaced basaltic sills at different stages of random-depth basalt emplacement (50 m sills every 10 kyr) for initial depth interval 10–15 km. Color-coded bars denote distance above sill margin (similar to those in Figures 7–10). Black is newly emplaced basalt, and gray is solidified mafic sill. Bolded temperature above each column marks the ambient crustal temperature at the time/depth of Xth sill emplacement (shown in adjacent parenthesis). The number at the base of the column marks the time that basalt sill center reaches 80% crystals, at which point exsolution of H₂O fluid is underway. Red and blue curves mark the thermal and melt fraction profiles (0.5 GPa; Figure 1c) along the first 50 m of wall rock adjacent to newly emplaced sills at the time indicated at the base of each column.

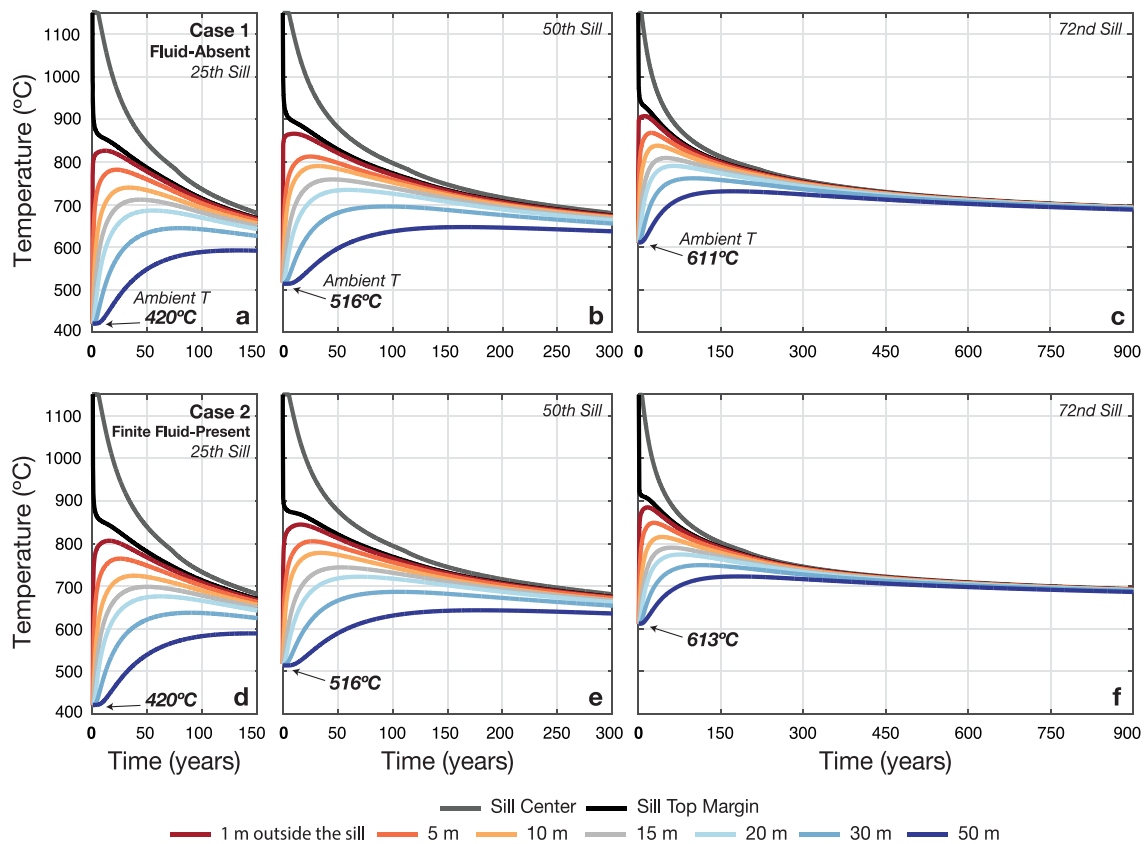


Figure 7. Temperature versus time immediately after sill emplacement at the basaltic sill center (gray line), basaltic sill top margin (black line), and depths 1, 5, 10, 15, 20, 30, and 50 m above the sill margin (colored lines same as in Figure 6). All cases show random emplacement of a 50-m thick sill every 10 kyr into an initial depth range of 10–15 km. (a) the 25th sill, (b) the 50th sill, and (c) the 72nd sill, all for fluid-absent F-T curves at 0.5 GPa (Figure 1b). (d–f) same as (a–c) but for finite fluid-present F-T curves at 0.5 GPa (Figure 1c). Bolded temperature with arrow is the ambient temperature of the crust immediately prior to sill emplacement.

fluid-present F-T curves, (4) sill thickness, and (5) the distribution of latent heat in basaltic sills across the liquidus-solidus interval.

4.2.1. Temperature Changes Within Hundreds of Years of Sill Emplacement

4.2.1.1. Heating of Adjacent 50 m of Wall Rock

After each 50-m basalt sill is emplaced, the ambient temperature of the crustal column steadily increases (Figure 2). Therefore, the temperature of the adjacent wall rock at the time of each sill emplacement varies with time (and location within the crustal column). In Figures 7 and 8, four different cases of temperature versus time for the newly emplaced sill and the adjacent 50 m of wall rock above the sill are shown (see Figures S8 and S9, for temperature vs. time plots at 0.3 GPa). Ambient temperature refers to the temperature of the wall rock immediately prior to sill emplacement. Each column of subfigures in Figure 7 refers to a different wall-rock lithology (Figure 6), which becomes more variable as sill emplacement progresses. The two rows in Figures 7 and 8 compare fluid-absent and finite fluid-present F-T curves (Figure 1) for the adjacent wall rock. The ambient temperature at the time each basaltic sill was emplaced is also noted, and there is some minor variation between fluid-absent and finite fluid-present scenarios.

In all cases (Figure 7), within the first 50 years of sill emplacement, there is an abrupt jump in temperature within the adjacent wall rock. When the ambient temperature is ~420 °C, temperatures reach ≥700 °C for 50 years within 15 m of adjacent wall rock. When the ambient temperature is ~516 °C, temperatures are ≥700 °C for 150 years within 20 m of adjacent wall rock. When the ambient temperature is ~611 °C, temperatures reach ≥700 °C for 500 years (and are ≥680 °C for 900 years) within 50 m of the sill margin. When the ambient temperature is ~657 °C (Figure 8), temperatures are ≥700 °C for >1,500 years (and are ≥680 °C for 3 kyr) within 50 m of the adjacent wall rock.

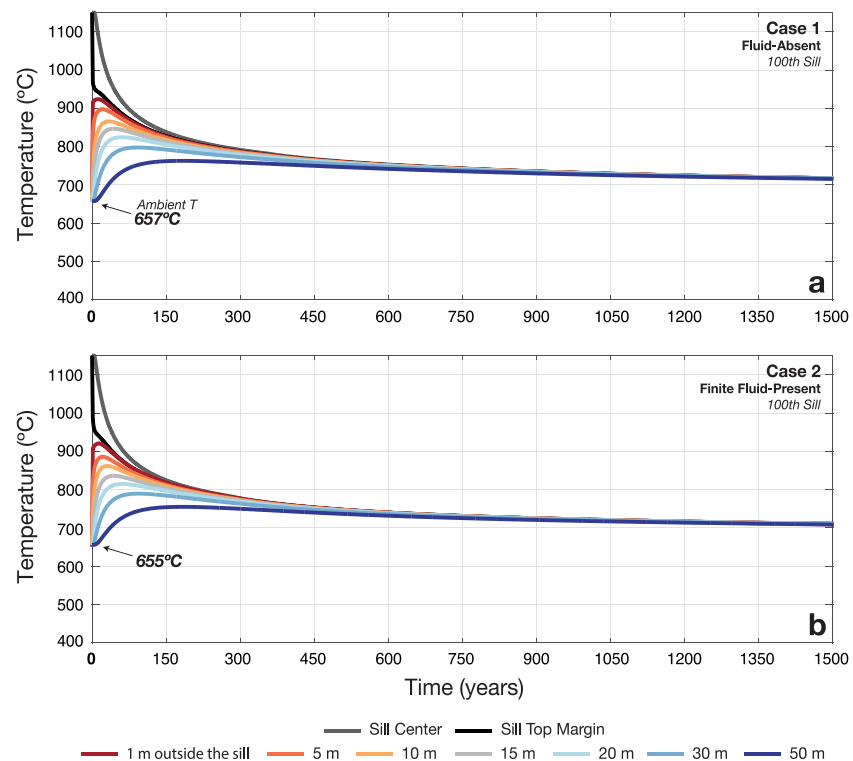


Figure 8. (a) Same as Figure 7a, but for the 100th sill; fluid-absent F-T curve. (b) Same as Figure 7d, but for the 100th sill; finite fluid-present F-T curve at 0.5 GPa. Bolded temperature with arrow is the ambient temperature of the crust immediately prior to sill emplacement.

Not surprisingly, the most important control on how high temperatures are reached in the wall rock adjacent to a newly emplaced sill, and for how long, is the ambient temperature. The effect of different wall rock lithologies on the thermal profile is minimal. The temperature profiles are slightly cooler when finite fluid-present versus fluid-absent F-T curves are employed, owing to the effect of latent heat of fusion during partial melting of wall rock (see section 4.1.7).

4.2.1.2. Cooling of Basaltic Sills

As ambient temperatures increase at the time of sill emplacement, the time it takes for a basaltic sill to cool to its solidus ($\sim 680^\circ\text{C}$ at 0.5 GPa; Figure 1a) also increases (Figure 7). When the ambient temperature is $\sim 420^\circ\text{C}$ versus 516°C , the basaltic sill center cools to 680°C within ~ 150 versus ~ 300 years. When the ambient temperature is $\sim 611^\circ\text{C}$ versus 657°C (Figure 7 and 8), it takes more than 900 versus 1,500 years for the sill center to reach its solidus, respectively.

4.2.2. Melt Fraction Changes Within Hundreds of Years of Sill Emplacement

4.2.2.1. Adjacent 50 m of Wall Rock

Melt fractions are calculated during each simulation alongside temperature, based on the various F-T curve relationships (Figure 1), within each wall rock lithology (granitoid vs. mafic sill). The amount of water that is present (i.e., fluid-absent vs. finite fluid-present conditions) plays a primary role in the extent of wall rock partial melting. Plots of melt fraction versus time since sill emplacement are shown in Figures 9 and 10 (see Figures S10 and S11, for melt fraction vs. time plots at 0.3 GPa). Each column of subfigures in Figures 9 and 10 refers to a different ambient temperature and wall-rock lithology (Figure 6), whereas the two rows compare fluid-absent and finite fluid-present F-T curves for the adjacent wall rock (Figure 1).

In all cases (Figure 9 and 10), within the first 50 years of sill emplacement, there is an abrupt jump in melt fraction within the adjacent wall rock. For the case of the 25th sill emplacement, the ambient temperature is $\sim 420^\circ\text{C}$ and the adjacent 50 m of wall rock is all granitoid. Here, the degree of melting in the adjacent wall rock ($\leq 20\%$ and $\leq 40\%$, respectively) is most strongly controlled by fluid-absent vs. finite fluid-present

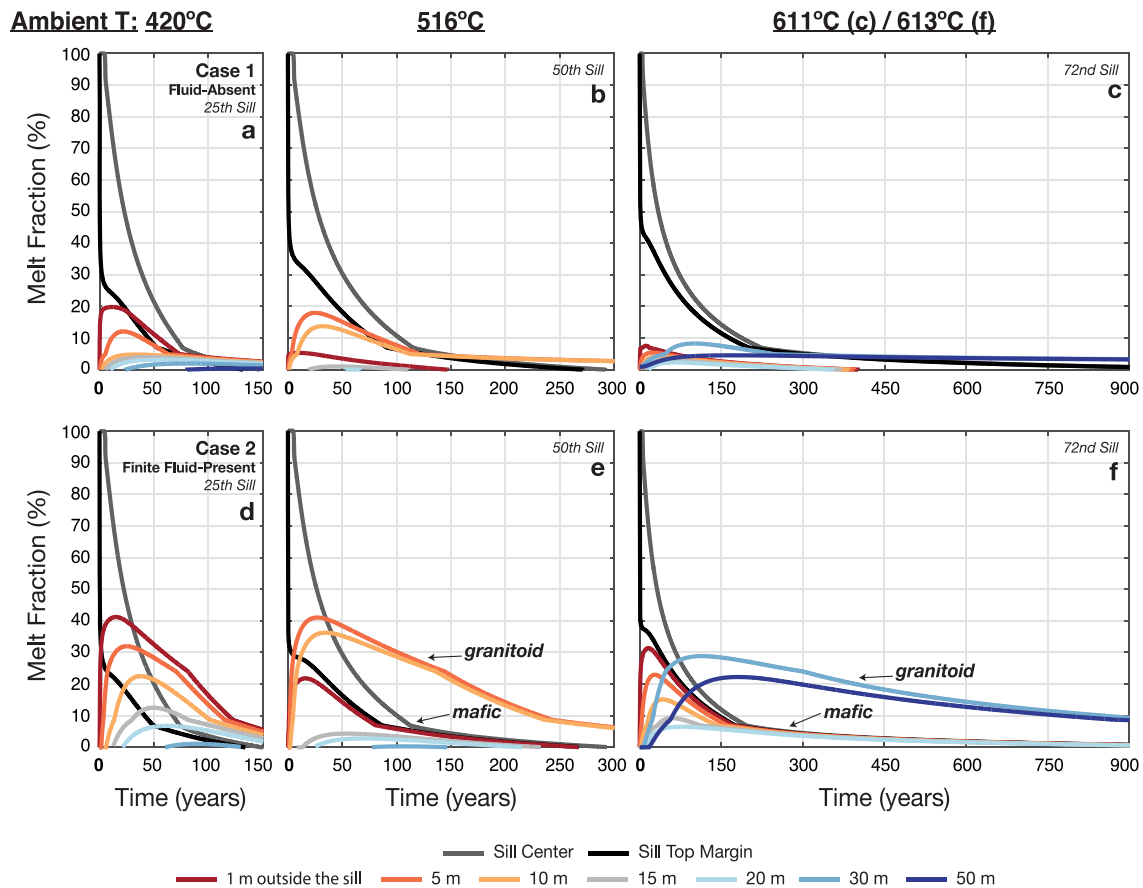


Figure 9. Same as Figure 7, but melt fraction instead of temperature is shown versus time immediately after sill emplacement. All symbols and colors are the same as in Figure 7. Arrows (in e–f) point out the different lithology (granitoid versus mafic sill) of the wall rock (see Figure 6).

conditions (Figure 9a and 9d). In the case of the 50th versus 72nd sill emplacements, an additional control on the degree of melting is the lithology of the wall rock adjacent to each newly emplaced sill. For example, owing to the random emplacement of basaltic sills, by chance the wall rock immediately above the 50th sill is predominantly mafic sill, with a small lens of granitoid 4–11 m above the sill margin, whereas the wall rock above the 72nd sill is mafic sill from 0–29 m and then granitoid (Figure 6). The higher melt fractions when the ambient temperature is lower (i.e., 516 vs. 611 °C) are due to the presence of granitoid within 10 m of the 50th sill's margin and its absence within the first 29 m of the 72nd sill. As expected, melt fractions are higher for finite fluid-present ($\leq 40\%$ melt) versus fluid-absent ($\leq 18\%$ melt) conditions in all cases.

In Figure 10, the melt fractions in the 50 m above the 100th sill, where the ambient temperature at the time of sill emplacement was ~ 657 °C, are shown for fluid-absent versus finite fluid-present conditions. In this case, the wall rock immediately above the sill is granitoid (0–19 m) and then solidified mafic sill (Figure 6). Under finite fluid-present conditions, the granitoid reaches 45–60% melt fraction and remains above 30% melt fraction for ~ 400 years, whereas the mafic sill reaches $\leq 12\%$ melt and remains above 5% melt for ~ 400 years. The temperature and melt fraction profile through the wall rock 50-m column is summarized in Figure 6 for various sill emplacements (at the point in time when the basaltic sill center has reached 80% crystallization).

4.2.2.2. Crystallization of Basaltic Sills

As ambient temperatures increase at the time of sill emplacement, the time it takes for the center of the basaltic sills to crystallize to 20% melt fraction also increases (Figure S7). It is at this stage of crystallization that the exsolution (and loss) of an H_2O -rich fluid phase from the basaltic sill begins, as discussed below. When the ambient temperature is 420 versus 516 °C, the basaltic sill center is 80% crystallized within ~ 50

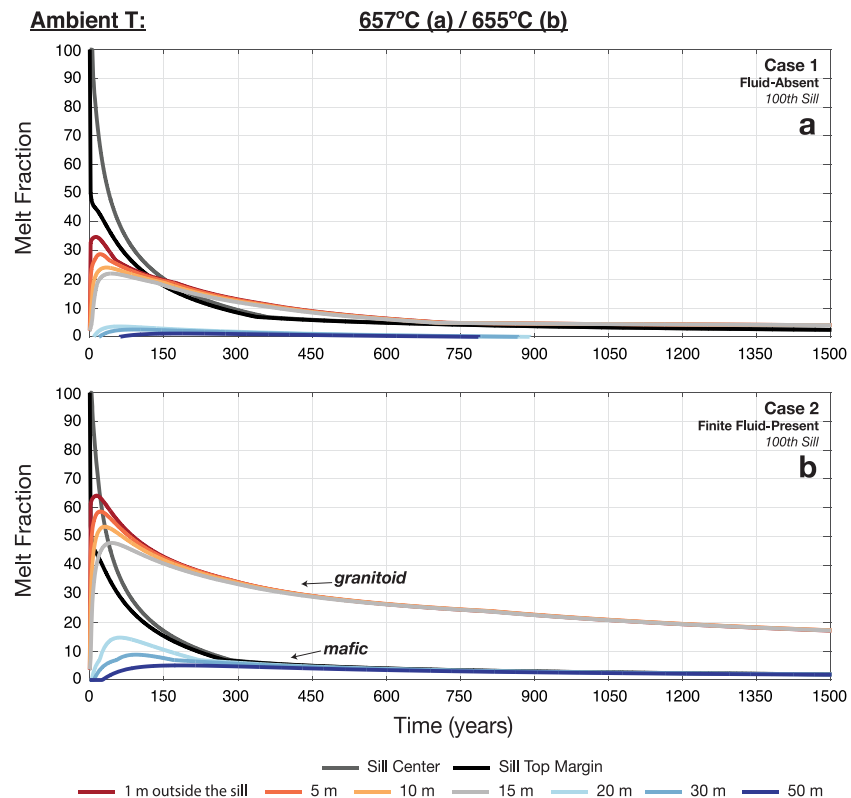


Figure 10. Same as Figure 8, but melt fraction instead of temperature is shown versus time immediately after sill emplacement. All symbols and colors are the same as in Figure 7. Arrows (in b) point out the different lithology (granitoid versus mafic sill) of the wall rock (see Figure 6).

versus ~69 years. When the ambient temperature is 611 and 657 °C (Figure 9 and 10), it takes ~101 and ~128 years, respectively, for the sill center to reach 80% crystallization.

4.2.2.3. Effect of Variable Thickness of Basaltic Sills

There are two causes for why varying the thickness of newly emplaced basaltic sills (e.g., 10, 50, and 100 m) controls the degree of partial melting that develops in the adjacent wall rock for any specified ambient crustal temperature. The first cause is the markedly different crustal lithologies (i.e., distribution of mafic sills within granitoid) that develop, dependent on whether sill thicknesses are 50 m (Figure 6) or alternatively either 10 or 100 m (Figure 11). Because mafic sills and granitoid have distinctly different F-T curves, resulting melt fractions will vary. The second reason is that sill thickness strongly controls the temperature evolution in the surrounding wall rock (e.g., Annen, 2017; Furlong et al., 1991). Variations in the thermal profile of the adjacent wall rock at specified ambient crustal temperatures for 10- and 100-m thick sills are shown in Figures S12–S15. Of relevance to this study is the impact of varying sill thickness on the extent of partial melting that occurs along wall-rock margins.

In the case of a 10-m thick sill emplaced into the crust at ~613 °C (Figure 12f), <10% melt fraction will form in granitoid wall rock (finite fluid-present) for ≤5 years within 10 m of the sill margin. In contrast, under similar conditions (ambient temperature and finite fluid-present), a newly emplaced 100-m thick sill causes ≥30% partial melting in granitoid wall rock for ~800 years up to a 30-m distance from the sill margin (Figure 13f). Moreover, when the ambient temperature is as low as ~431 °C, the emplacement of a 100-m thick sill induces significant partial melting (≥20%) of granitoid (finite fluid-present conditions) for ≤300 years within 20 m from the sill margin (Figure 13d). Importantly, the results in Figure 12 (10-m sills) and Figure 13 (100-m sills) are both markedly different from those obtained under similar conditions for 50-m thick basaltic sills (Figure 10) (see Figures S16 and S17 for melt fraction results for 10- and 100-m thick sills when ambient temperature is ~650 °C). The key observation is that variations in the thickness of

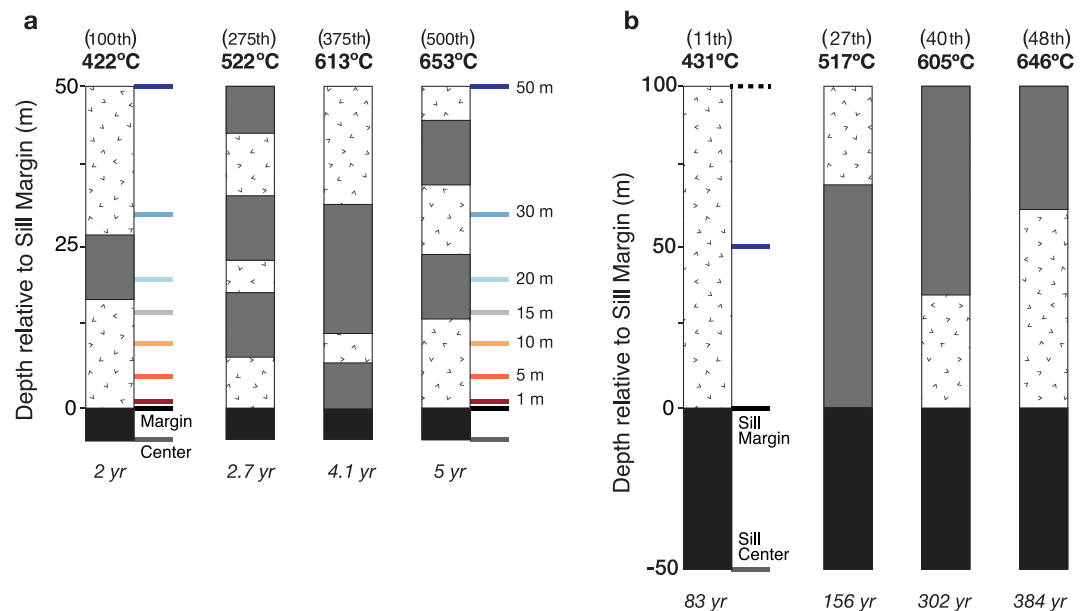


Figure 11. Examples of a resultant lithologic column (a) 0–50 m above newly emplaced 10-m thick basaltic sills (10 m/2 kyr) and (b) 0–100 m above newly emplaced 100-m thick basaltic sills (100 m/20 kyr) at different stages of emplacement. The initial depth interval is 10–15 km in both cases. Color-coded bars denote distance above sill margin (similar to that in Figure 6). Black is newly emplaced basalt, and gray is solidified mafic sill. Bolded temperature above each column marks the ambient crustal temperature at the time/depth of Xth sill emplacement (shown in adjacent parenthesis). The number at the base of the column marks the time that basalt sill reaches 80% crystallization, at which point exsolution of H_2O fluid is underway. Sill scale results for (a) are shown in Figures 12, S12, S13, and S16. Sill scale results for (b) are shown in Figures 13, S14, S15, and S17.

newly-emplaced sills (10, 50, and 100 m) lead to significant differences in the degree and duration of wall-rock partial melting.

4.2.3. Effect of a Variable Distribution of Latent Heat Across the Liquidus-Solidus Interval

An additional factor that may affect the cooling history of a basaltic sill and the transient heating of the adjacent wall rock is a nonlinear distribution of latent heat across the liquidus-solidus interval. For the cases shown in Figures 7–10, latent heat is distributed linearly. In the two alternative cases (Figure S4), where (1) 75% of the latent heat is released in the first 25% of crystallization and (2) 75% of the latent heat is released in the last 25% of crystallization, the time it takes the 100th basaltic sill to cool to 810 °C (~80% crystallization at 0.3 GPa; Figure 1a), is ~6 years longer and ~29 years shorter, respectively, compared to the ~120 years it takes for a linear distribution of latent heat (Figure S4c). Similarly, the time interval over which the adjacent wall rock, 10 m from sill margin, remains ≥ 740 °C ($\geq 30\%$ melt fraction; finite fluid-present conditions, Figure 1c) is ~9 years longer and ~37 years shorter, respectively, for Cases a and 2 (Figure S4a).

The most important parameter to track is the difference in time between the duration of wall rock heating (e.g., 10 m above sill ≥ 740 °C) and the time it takes the basaltic sill to crystallize to 80% (i.e., to 810 °C; Figure 1a), which marks the onset of fluid exsolution (see section 4.3). For a linear distribution of latent heat, this time interval for the 100th sill (Figure 6) is 275 years, whereas for Cases 1 and 2 (Figure S4a), the time interval is 3 years longer and 8 years shorter, respectively. Therefore, a nonlinear distribution of latent heat has little effect on the relative time scale between the cooling of a basaltic sill and the transient heating of the adjacent wall rock.

4.3. Transfer of Water from Basaltic Sills to Partial Melt of Adjacent Wall Rock

A final process we explore is the transfer of dissolved water from the basaltic sills into partial melts of the adjacent wall rock and what conditions are required for this to occur. For example, the solubility limit for dissolved water in the melt is ~10 wt% in a basalt that is emplaced at a depth with an ambient pressure of

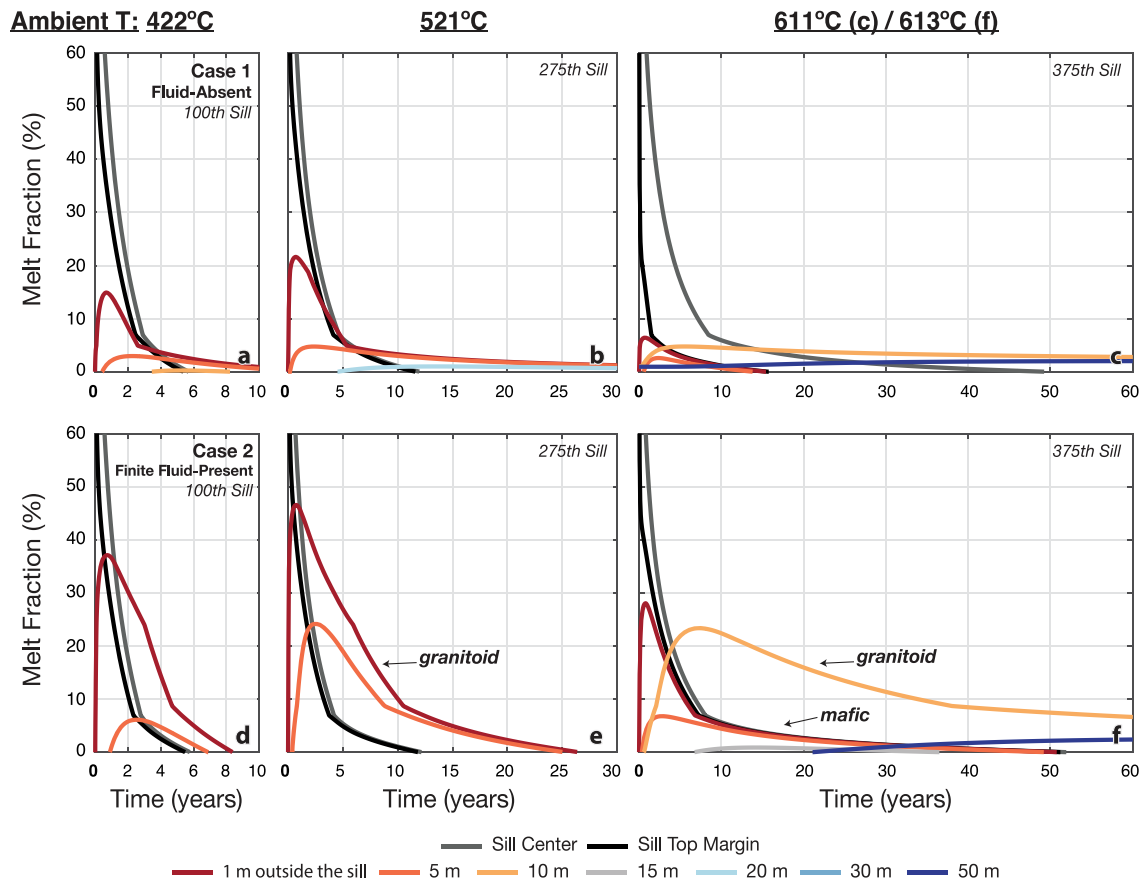


Figure 12. Melt fraction versus time immediately after 10-m thick sill emplacement at the basaltic sill center (gray line), basaltic sill top margin (black line), and depths 1, 5, 10, 15, 20, 30, and 50 m above the sill margin (colored lines the same as in Figure 11). All cases are random emplacement of 10-m thick sill every 2 kyr into an initial depth range of 10–15 km. (a) the 100th sill, (b) the 275th sill, and (c) the 375th sill from a single simulation realization, all for fluid-absent F-T curves at 0.5 GPa (Figure 1b). (d–f) same as (a–c) but for finite fluid-present F-T curves at 0.5 GPa (Figure 1c). Arrows (in e–f) point out different lithology (granitoid vs. mafic sill) of the wall rock (see Figure 11).

0.5 GPa (e.g., Behrens et al., 2001; Zhang et al., 2007). If the initial water content is ~3 wt% (Gazel et al., 2012), 10 wt% H₂O will be reached in the melt during late-stage crystallization. Therefore, exsolution of an H₂O-rich fluid phase is expected to occur by the time the sill center reaches 80% crystallization. Based on results from Bachmann and Bergantz (2006), migration of an exsolved fluid phase (via “gas sparging”) through partially molten sill and wall rock is expected, and transport distances of 5–10 m are plausible in ~5–40 years. The question is whether the time duration for partial melting in the adjacent wall rock (e.g., Figure 14) matches the time scale required for the basaltic sill to crystallize ~80%, enabling transfer of a fluid phase to the adjacent wall rock.

In Figures 14a and 14b, the time it takes (relative to the instant of sill emplacement) for granitoid wall rock at specified distances (1, 5, 10, 20, 30, and 50 m) above the sill margin to reach a certain temperature, and thus melt fraction (finite fluid-present conditions), is shown. Superimposed on these plots are isopleths of melt fraction (10% and 20%) at each depth, based only upon the sill emplacements where all of the material up to that depth was granitoid. The valley-shaped pattern for an isopleth of melt fraction in the wall rock, at a given distance from the sill margin, illustrates the initial, transient heating of the wall rock (i.e., at a given location relative to the sill margin, 10% melt fraction will be achieved earlier as ambient temperature increases), followed by the subsequent cooling of the wall rock (at a given location relative to the sill margin, as ambient temperature increases, it takes longer to dip below 10% melt fraction again). Pairs of same-colored points at the same ambient temperature denote the times after a given sill emplacement in which the wall rock partially melts above or solidifies below a certain melt fraction, where dot color corresponds to the distance from the sill margin. Spacing between those pairs of points is the time interval in

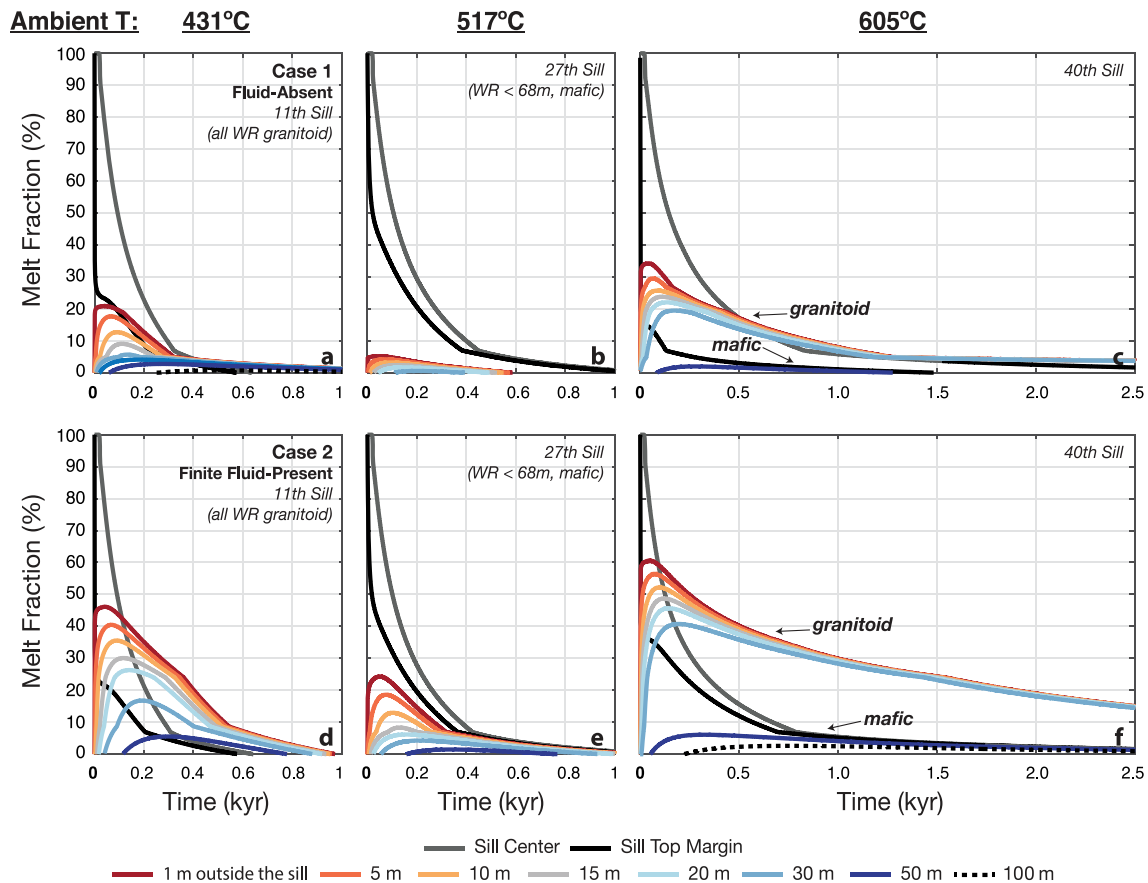


Figure 13. Melt fraction versus time immediately after 100-m sill emplacement at the basaltic sill center (gray line), basaltic sill top margin (black line), and depths 1, 5, 10, 15, 20, 30, 50, and 100 m above the sill margin (colored lines the same as in Figure 11). All cases are random emplacement of 100-m thick sill every 20 kyr into an initial depth range of 10–15 km. (a) the 11th sill, (b) the 27th sill, and (c) the 40th sill from a single simulation, all for fluid-absent F-T curves at 0.5 GPa (Figure 1b). (d–f) same as (a–c) but for finite fluid-present F-T curves at 0.5 GPa (Figure 1c). Arrows (in c and f) point out different lithology (granitoid vs. mafic sill) of the wall rock (see Figure 11).

which the wall rock at that distance from the sill has melt fraction above the indicated level. Also shown in this plot is a solid red curve that marks the time at which the center of a recently emplaced basaltic sill reaches 80% crystallization (Figure S7b) and therefore exsolution of an H_2O -rich fluid phase. Tie lines located to the right of this curve mark the time interval where fluid-present conditions in granitoid wall rock are most likely to exist, owing to transfer of an exsolved H_2O -rich fluid. Therefore, the relevant melt fraction curves for both types of wall rock, granitoid, and solidified mafic sills are those for finite fluid-present conditions (Figure 1c), which are also used for the melt fractions illustrated in Figure 6.

In Figures 14c and 14d, a similar pair of plots is shown, but here, the wall rock is solidified mafic sills and the isopleths are for 5% and 10% melt fractions (calculated for finite fluid-present conditions; Figure 1c) at each of the depths considered. The isopleths are based only upon the sill emplacements where all of the wall rock up to that depth was solidified mafic sill. The more refractory nature (i.e., lower melt fraction at a given T and H_2O content) of the mafic sills relative to granitoid is illustrated in Figures 14a–14d. Melt fractions of ~5–10% do not develop under finite fluid-present conditions until ambient temperatures are $\geq 500^\circ C$ and only within 1–5 m of the newly emplaced sill margin. In sharp contrast, melt fractions in finite fluid-present granitoid up to 20 m from the sill margin exceed 20% for ~20 years and 10% for ~100 years. Despite the relatively refractory nature of the solidified mafic sills, their close spatial proximity to granitoid increases the chance that small degree partial melts from the mafic sills will develop hydraulic connectivity with the larger degree partial melts from the granitoid, leading to interaction and mixing of compositionally distinct partial melts derived from two very different source rocks.

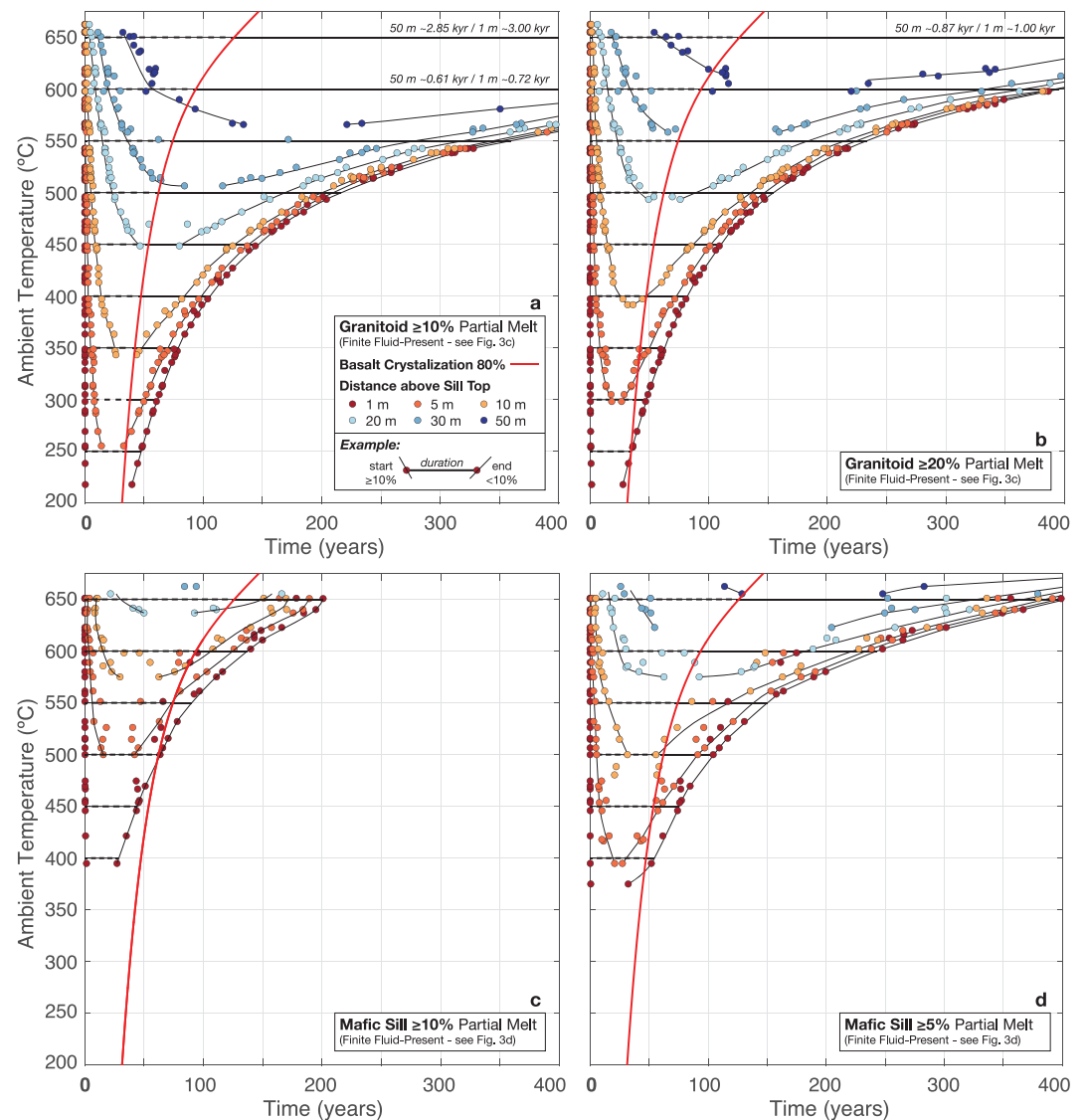


Figure 14. (a) Plot of ambient temperature versus time with superimposed isopleths of finite fluid-present granodiorite at fixed distances above the basaltic sill margin (color-coded; see legend) that contain $\geq 10\%$ melt fraction at 0.5 GPa. Red line is the time it takes a newly emplaced basaltic sill to crystallize 80% (onset of fluid exsolution; Figure S7). Bold isothermal lines denote time interval, for a specific initial ambient temperature, over which granitoid wall rock is under finite fluid-present conditions and thus contains $\geq 10\%$ melt fraction. (b) Same as (a) but for granitoid with $\geq 20\%$ melt fraction. (c) Same as (a) but for solidified mafic sills with $\geq 10\%$ melt fraction. (d) Same as (a) but for solidified mafic sills with $\geq 5\%$ melt fraction. Isoleth curves that connect the color-coded dots are hand drawn.

The efficient transfer of H_2O fluid from a crystallizing sill to the wall rock is more clearly demonstrated when the loss (i.e., advection) of H_2O fluid, and thus heat, from the hydrous basalt during late-stage crystallization is accounted for. For example, the loss of $\sim 2 \text{ wt\%}$ H_2O fluid, with a heat content of $\sim 3,000 \text{ J/g}$ at 700°C and 0.5 GPa (Burnham et al., 1969) out of a basaltic sill ($> 80\%$ crystallized; $C_p = \sim 1.4 \text{ J/g-K}$; Table 1) will lead to $\sim 43^\circ\text{C}$ of cooling. Moreover, the addition of this H_2O fluid into 50 m of wall rock will similarly induce another $\sim 43^\circ\text{C}$ of heating and thus prolong the duration of the transient wall rock heating.

The conclusion, therefore, is that finite fluid-present F-T curves are relevant for granitoid and mafic sill lithologies in wall rock adjacent to newly emplaced sills. Once the ambient temperature is $\sim 600^\circ\text{C}$, significant melt fractions develop that contain dissolved H_2O (Figure 6) but at levels below solubility limits owing

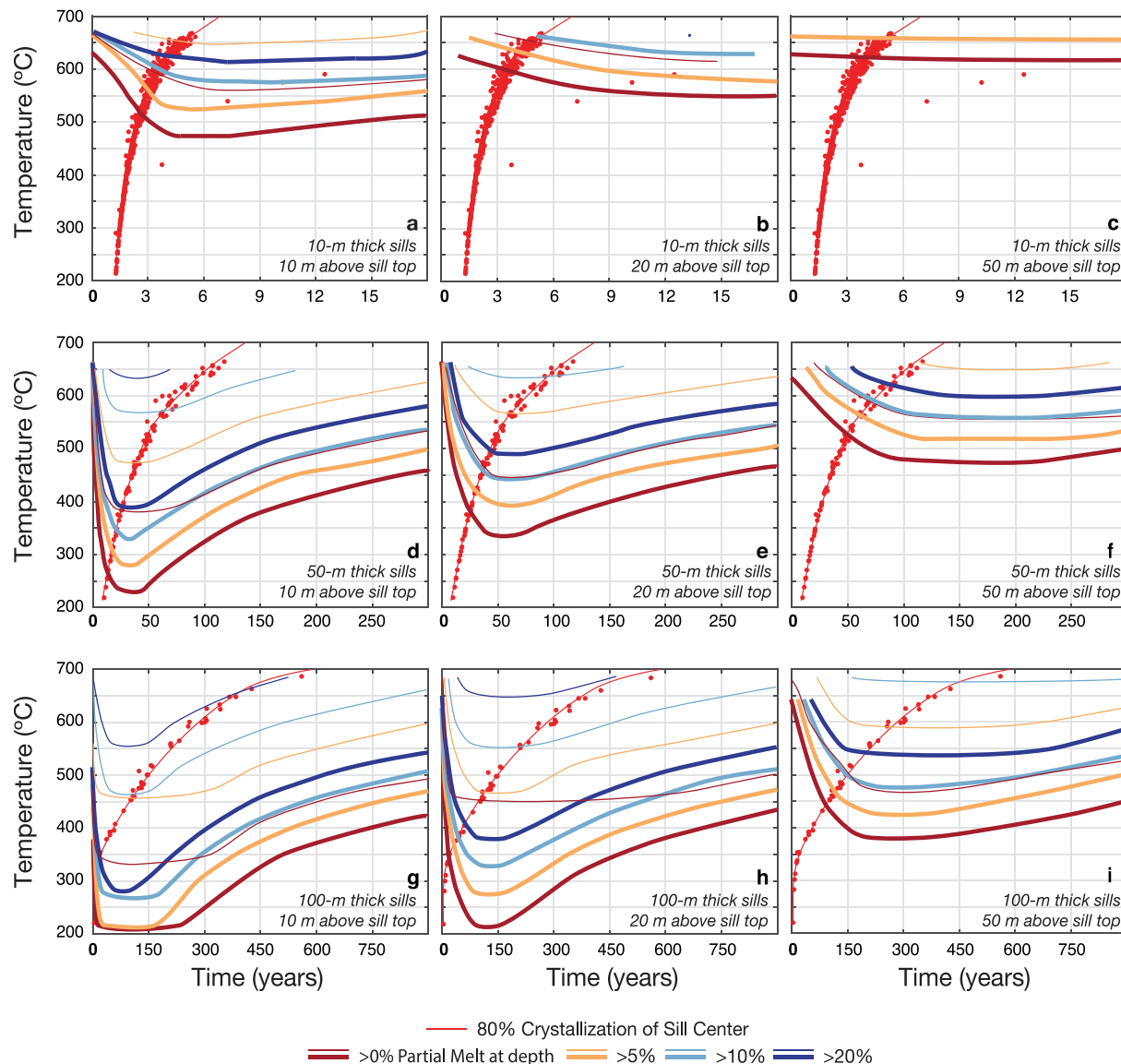


Figure 15. Plots of ambient temperature versus time with superimposed isopleths of finite fluid-present granodiorite (bold) and mafic sill (thin) partial melting windows. (a–c) Results for 10-m thick sills at 10, 20, and 50 m above the sill margin. (d–f) Results for 50-m thick sills at 10, 20, and 50 m above the sill margin. (g–i) Results for 100-m thick sills at 10, 20, and 50 m above the sill margin.

to the limited amount of H_2O available from the basaltic sill. The formation of hydrous partial melts that are fluid undersaturated has implications for their rapid segregation and transport as discussed below.

4.3.1. The Role of Variable Sill Thickness

Figure 15 presents a similar set of plots as shown in Figure 14, but in this case, the effect of varying sill thickness (10, 50, and 100 m) is featured. The comparative results clearly illustrate that partial melting along the wall-rock margins of newly emplaced sills is far more extensive for thicker (100-m) versus thinner (10-m) sills. In the case of 10-m thick sills, very limited and transient (<20 years) partial melts form in the wall-rock margin above newly emplaced sills, even when ambient temperatures are relatively high ($\geq 600^\circ\text{C}$). In contrast, in the case of 100-m thick sills, extensive (>20%) and long-lived (>300 years) melt fractions develop along a new sill margin, even when ambient temperatures are relatively low ($\sim 400^\circ\text{C}$). These results underscore the role of variable sill thickness in controlling the extent and duration of partial melting along the wall-rock margins of a newly emplaced sill, as well as the ambient temperature at which significant wall-rock partial melting can occur.

5. Discussion

5.1. Development of Crustal-Scale Thermal Profiles and Mixed Lithologies

Several factors influence, to varying degrees, the temperature-depth profile through the crustal column due to the influx of basaltic sills. Among the parameters explored in this study, three exert the strongest control: (1) random versus fixed emplacement of sills, (2) the initial depth range of sill emplacement, and (3) the time interval over which a thickness of basalt is emplaced. These results are broadly consistent with previous work in the literature (e.g., Annen et al., 2006; Colón et al., 2018; Dufek & Bergantz, 2005; Karakas & Dufek, 2015; Leeman et al., 2008). An important factor additionally explored in this study is the development of a mixed crustal lithology as the crust is gradually heated through the periodic and random emplacement of basaltic sills. Once temperatures reach values that facilitate partial melting of wall rock, both the original granitoid as well as previously solidified mafic sills are eligible source rocks, leading to two sets of compositionally distinct partial melts. Therefore, tracking the development of the mixed crustal lithology is as important as tracing the evolving thermal profile through the crust.

The results from this study show that crustal temperatures in the middle-upper crust can reach $>600^{\circ}\text{C}$ with varying proportions of mafic sills embedded within the granitoid crust, dependent on the rate of basalt emplacement, as well as the depth interval over which basalts are initially emplaced. For example, the higher the rate of emplacement, the less basalt is required to heat the crust to $>600^{\circ}\text{C}$ (Figure 4b) and thus the lower the ratio of solidified mafic sills to granitoid in the crustal column. Similarly, the deeper and broader the initial depth interval of basalt emplacement, the lower the basalt:granitoid ratio required to heat the surrounding crust to $>600^{\circ}\text{C}$ (Figure S6).

5.2. Importance of Hydrous Partial Melts of Wall Rock and Their Transport

An additional factor of potential relevance to crustal-scale thermal profiles is the formation, and then advection, of partial melts along the wall-rock margins of newly emplaced sills. The high spatial and temporal resolution of our numerical model reveals the thermal consequences in adjacent wall rock in response to subsequent injections of 50-m thick basaltic sills once ambient crustal temperatures reach $\sim 600^{\circ}\text{C}$. After a new sill is emplaced, temperatures in the adjacent wall rock will reach $\geq 700^{\circ}\text{C}$ for ≥ 500 years within 50 m of the sill margin. It is further shown that basaltic sills with an initial H_2O content of 3 wt% will undergo 80% crystallization at the sill center within 100 years when the ambient temperature of the surrounding crust is $\sim 600^{\circ}\text{C}$. After this much crystallization, the basaltic sill is expected to exsolve an H_2O -rich fluid phase, leading to finite fluid-present conditions for the adjacent wall rock at its solidus. Resulting melt fractions in granitoid and mafic sill wall rock will reach $\geq 20\%$ and $\geq 5\%$, respectively, within 30 and 20 m of the sill margin for ~ 300 and ~ 150 years, respectively. Given the close spatial association of mafic sills and granitoid in the crustal column (Figure 6), it is expected that the relatively low-degree ($\sim 5\%$) partial melts from the mafic sills will interact with relatively high-degree ($\sim 20\%$) partial melts from the granitoid wall rock.

Once ambient crustal temperatures reach $\sim 650^{\circ}\text{C}$, any subsequent emplacement of a 50-m sill will induce substantial partial melting of the adjacent wall rock (Figures 6 and 14). At relatively high melt fractions, the partial melts will be hydrous but are expected to be fluid undersaturated owing to the limited amount of water available from the basaltic sills. Under these conditions, a large volume increase is created during partial melting. At 500 MPa, any dissolved water content less than 10 wt% will be fluid undersaturated, and the density of a minimum/eutectic melt partial melt from a granitoid (with $\sim 5\text{--}9$ wt% dissolved H_2O) will range from $\sim 2,360\text{--}2,290\text{ kg/m}^3$ (Lange, 1997; Ochs & Lange, 1999), which represents a volume increase of $\sim 11\text{--}14\%$. Experimental studies confirm that partial melting reactions of crustal rocks that produce a positive change in volume lead to microcracking and melt-filled fractures (e.g., Connolly et al., 1997; Rushmer, 2001). The formation of near-vertical melt-filled fractures in granitoid wall rock will be further facilitated by the preexistence of scattered aplite dikes (Lange, 2019), which are ubiquitous in the Sierra Nevada batholith. The near-eutectic bulk composition (high- SiO_2 rhyolite) of aplite dikes ensures extensive melting (limited only by the finite amount of H_2O present), and they are likely sites for melt-filled fractures to readily develop.

The rapid development of melt-filled fractures (through melting of aplite dikes) will facilitate the extraction of partial melt from the wall rock (e.g., Sleep, 1988). Once the melt-filled fractures exceed the

critical length of a self-propagating dike, it is possible for segregation to occur through dike transport. The critical length at which a vertically oriented dike of constant volume can “self-propagate” is given by the relation $(K_c/[\Delta\rho g])^{2/3}$, where K_c is the fracture toughness ($\sim 1 \text{ MPa m}^{1/2}$; Gonnermann & Taisne, 2015). For a $\Delta\rho$ (rock-melt) of $\sim 290\text{--}360 \text{ kg/m}^3$, critical lengths of $\sim 43\text{--}50 \text{ m}$ are obtained, which are on the same length scale of the partially melted wall-rock margins. If these partial melts segregate under fluid-undersaturated conditions, their liquidus temperatures will decrease during ascent (owing to the positive dT/dP slope of their liquidus), which will prolong their transit through the crust despite conductive loss of heat (Waters & Lange, 2017).

The key point for this study is that rapid segregation and transport of these partial melts out of the wall rock is a viable hypothesis that merits further examination, because it will have major consequences to thermal profiles through the crust. The removal of these partial melts by advection will cool their source regions and heat shallower levels of the crust where they are emplaced. Moreover, this process of melt advection will also alter the composition of the middle-upper crust.

5.3. Compositional Reworking of Crust due to Transport of Wall-Rock Partial Melts

Over time, the episodic vertical transport of small volumes of rhyolitic melt (i.e., partial melts of wall rock) into higher levels of the crust will gradually lead to a relatively large accumulation of Plio-Quaternary leucogranite in the upper crust, with a composition that is more evolved than the Mesozoic granodiorite that characterizes most of the exposed granitoid in this region (e.g., Oliver et al., 1993). The rhyolitic partial melts that form along the margins of newly emplaced basalt sills will be derived from both previously solidified mafic sills and granitoid. Therefore, the two sets of partial melts will be characterized by distinctly different major- and trace-element concentrations, and these two distinctly different partial melts are likely to interact during their segregation and transport, given the close spatial proximity of solidified mafic sills within granitoid wall rock (Figure 6). It is therefore likely that Plio-Quaternary leucogranite bodies of rhyolitic composition (73–75 wt% SiO_2) will form in the upper crust with major-element compositions dominated by the more voluminous partial melts of granitoid, but with concentrations of select trace elements (e.g., Sr) that are predominantly derived from the Sr-rich mafic sills (Cousens, 1996). Subsequent partial melting of these leucogranite bodies will lead to high- SiO_2 (76–77 wt%) rhyolites with relatively low Sr concentrations but Sr-isotopic values that mirror those of the associated Long Valley basalts (Lange, 2013), which is a key feature of the Long Valley rhyolites (e.g., Davies et al., 1994; Halliday et al., 1984; Hildreth & Wilson, 2007; Metz & Mahood, 1991).

6. Conclusions and Future Work

There are three broad sets of conclusions that can be drawn from this study regarding the thermal evolution of the crust beneath the Long Valley caldera in eastern California, in response to the emplacement of basaltic magma. First, on the scale of tens of kilometers, as the crustal column gradually heats up due to the random emplacement of basaltic sills, a mixed crustal lithology (solidified mafic sills as well as granitoid) will develop. Importantly, the amount of basalt required to heat $>10 \text{ km}$ of the crustal column to $\sim 600^\circ\text{C}$ varies strongly with basalt emplacement rate and the depth interval of basalt emplacement. Therefore, the proportion of mafic sills to granitoid in the crustal column can vary by more than a factor of two (e.g., $\sim 1:1$ to $<1:2$), once crustal temperatures reach $\sim 600^\circ\text{C}$. The abundance and composition of subsequent partial melts of this mixed crustal lithology, at a given $T\text{-}P_{\text{H}_2\text{O}}$ condition, will depend strongly on the ratio of mafic sill to granitoid. Therefore, tracking the development of the mixed crustal lithology is as important as tracing the evolving thermal profile through the crust.

A second major conclusion is that in order to accurately track the thermal evolution of the crust due to the influx of basaltic sills, attention must be paid to the transient heating and partial melting of the wall rock adjacent to newly emplaced sills, which depends on the thickness of the basaltic sills as well as the time scale and mechanism by which water is transferred from them to the adjacent wall rock. This requires numerical modeling at high spatial (1 m) and temporal (tens to hundreds of years) resolution, as shown in Figures 7–10. The time scale for the crystallization of the basaltic sills (hundreds of years or less), resulting in the exsolution of a fluid phase, matches the time interval of transient heating of the wall rock. This indicates an effective transfer of water from the crystallizing sill to the wall rock, therefore finite fluid-present conditions

prevail and extensive fluid-undersaturated partial melting of wall rock is expected, even when ambient temperatures are below the water-saturated solidus. Moreover, the wall-rock partial melts are expected to be derived from both granitoid and previously solidified mafic sills, and the two sets of compositionally distinct melts are likely to interact during segregation and ascent.

A third conclusion is that some of these wall-rock partial melts may undergo rapid segregation and transport via fractures, due to large volume increases during partial melting and the development of melt-filled fractures within the adjacent wall rock (owing to melting of preexisting aplite dikes), which exceed the critical length for dike self-propagation. The advection of these partial melts to higher levels in the crust via fractures will strongly affect the thermal and compositional profile through the middle and upper crust. This is an area of future work, and we are currently developing a new numerical 2-D model that will include the advection of these partial melts, operating at the same high spatial (1 m) and temporal (1 hr) resolution as the 1-D model presented in this study.

Acknowledgments

This research was supported by the National Science Foundation (EAR-1855751), as well as discretionary research funds from the University of Michigan (UM). The first author was partially supported by a graduate student fellowship from the Department of Earth and Environmental Sciences at the UM. The model code, a run file, and representative output can be found at DOI: 10.5281/zenodo.3549797; more output available upon request.

References

- Andersen, N. L., Jicha, B. R., Singer, B. S., & Hildreth, W. (2017). Incremental heating of Bishop Tuff sanidine reveals preruptive radiogenic Ar and rapid remobilization from cold storage. *Proceedings of the National Academy of Sciences of the United States of America*, 114(47), 12,407–12,412. <https://doi.org/10.1073/pnas.1709581114>
- Anderson, A. T., Davis, A. M., & Lu, F. (2000). Evolution of Bishop Tuff rhyolitic magma based on melt and magnetite inclusions and zoned phenocrysts. *Journal of Petrology*, 41(3), 449–473. <https://doi.org/10.1093/petrology/41.3.449>
- Annen, C. (2017). Factors affecting the thickness of thermal aureoles. *Frontiers in Earth Science*, 5, 82. <https://doi.org/10.3389/feart.2017.00082>
- Annen, C., Blundy, J. D., & Sparks, R. S. J. (2006). The genesis of intermediate silicic magmas in deep crustal hot zones. *Journal of Petrology*, 47(3), 505–539. <http://doi.org/10.1093/petrology/egi084>
- Bachmann, O., & Bergantz, G. W. (2006). Gas percolation in upper-crustal silicic crystal mushes as a mechanism for upward heat advection and rejuvenation of near-solidus magma bodies. *Journal of Volcanology and Geothermal Research*, 149, 85–102. <https://doi.org/10.1016/j.jvolgeores.2005.06.002>
- Bailey, R. A. (1989). Geologic map of the Long Valley caldera, Mono-Inyo craters volcanic chain, and vicinity, eastern California. U.S. Geological Survey Map I-1933. <https://doi.org/10.3133/i1933>
- Bateman, P. C. (1992). Plutonism in the central part of the Sierra Nevada batholith, California. *U.S. Geological Survey Professional Paper*, 1483, 1–186.
- Behrens, H., Meyer, M., Holtz, F., Benne, D., & Nowak, M. (2001). The effect of alkali ionic radius, temperature, and pressure on solubility of water in MAlSi_3O_8 melts (M = Li, Na, K, Rb). *Chemical Geology*, 174, 275–289. [http://doi.org/10.1016/S0009-2541\(00\)00320-X](http://doi.org/10.1016/S0009-2541(00)00320-X)
- Bindeman, I. N., & Simakin, A. G. (2014). Rhyolites—Hard to produce, but easy to recycle and sequester: Integrating microgeochemical observations and numerical models. *Geosphere*, 10(5), 930–957. <https://doi.org/10.1130/GES00969.1>
- Bohrson, W. A., & Spera, F. J. (2001). Energy-constrained open system magmatic processes II: Application of energy-constrained assimilation–fractional crystallization (EC-AFC) model to magmatic systems. *Journal of Petrology*, 42, 1019–1041. <https://doi.org/10.1093/petrology/42.5.1019>
- Burnham, C. W., Holloway, J. R., & Davis, N. F. (1969). Thermodynamic properties of water to 1000 °C and 10,000 bars. Research and Development Progress Report No. 414, US Department of the Interior. <https://doi.org/10.1130/SPE132>
- Chapman, D. S., & Furlong, K. P. (1992). Thermal state of the continental lower crust. In D. M. Fountain, R. Arculus, & R. W. Kay (Eds.), *Continental lower crust. Developments in geotectonics* (Vol. 23, pp. 179–199). Amsterdam: Elsevier.
- Christiansen, R. L. (2001). The Quaternary and Pliocene Yellowstone volcanic field of Wyoming, Idaho, and Montana. *US Geological Survey Professional Paper*, 729(G), G1–G145. <https://doi.org/10.3133/pp729G>
- Coleman, D. S., Bartley, J. M., Glazner, A. F., & Pardue, M. J. (2012). Is chemical zonation in plutonic rocks driven by changes in source magma composition or shallow-crustal differentiation? *Geosphere*, 8(6), 1568–1587.
- Colón, D. P., Bindeman, I. N., & Gerya, T. V. (2018). Thermomechanical modeling of the formation of a multilevel, crustal-scale magmatic system by the Yellowstone plume. *Geophysical Research Letters*, 45, 3873–3879. <https://doi.org/10.1029/2018GL077090>
- Connolly, J. A. D., Holness, M. B., Rubie, D., & Rushmer, T. (1997). Reaction-induced microcracking: An experimental investigation of a mechanism for enhancing anatectic melt extraction. *Geology*, 25, 591–594. [http://doi.org/10.1130/0091-7613\(1997\)025<0591:RIMAEI>2.3.CO](http://doi.org/10.1130/0091-7613(1997)025<0591:RIMAEI>2.3.CO)
- Cousens, B. L. (1996). Magmatic evolution of Quaternary mafic magmas at Long Valley caldera and the Devils Postpile, California: Effects of crustal contamination on lithospheric mantle-derived magmas. *Journal of Geophysical Research*, 101(B12), 27,673–27,689. <https://doi.org/10.1029/96JB02093>
- Davies, G. R., Halliday, A. N., Mahood, G. A., & Hall, C. M. (1994). Isotopic constraints on the production rates, crystallization histories and residence times of pre-caldera silicic magmas, Long Valley, California. *Earth and Planetary Science Letters*, 125(4), 17–37. [https://doi.org/10.1016/0012-821X\(94\)90204-6](https://doi.org/10.1016/0012-821X(94)90204-6)
- Dormand, J. (1996). *Numerical methods for differential equations: A computational approach* (p. 368). New York: CRC Press LLC.
- Du Bray, E. A., John, D. A., Cousens, B. L., Hayden, L. A., & Vikre, P. G. (2016). Geochemistry, petrologic evolution, and ore deposits of the Miocene Bodie Hills volcanic field, California and Nevada. *American Mineralogist*, 101(3), 644–677. <https://doi.org/10.2138/am-2016-5440>
- Dufek, J., & Bergantz, G. W. (2005). Lower crustal magma genesis and preservation: A stochastic framework for the evaluation of basalt–crust interaction. *Journal of Petrology*, 46(11), 2167–2195. <https://doi.org/10.1093/petrology/egi049>
- Flügelner, M. M., Klempner, S. L., & Christensen, N. I. (2000). Three-dimensional seismic model of the Sierra Nevada arc, California, and its implications for crustal and upper mantle composition. *Journal of Geophysical Research*, 105(B5), 10,899–10,922. <https://doi.org/10.1029/2000JB900029>

- Flinders, A. F., Shelly, D. R., Dawson, P. B., Hill, D. P., Tripoli, B., & Shen, Y. (2018). Seismic evidence for significant melt beneath the Long Valley Caldera, California, USA. *Geology*, 46(9), 799–802.
- Frassetto, A. M., Zandt, G., Gilbert, H., Owens, T. J., & Jones, C. H. (2011). Structure of the Sierra Nevada from receiver functions and implications for lithospheric foundering. *Geosphere*, 7(4), 898–921. <https://doi.org/10.1130/GES00570.1>
- Furlong, K. P., Hanson, R. B., & Bowers, J. R. (1991). Modeling thermal regimes. In D. M. Kerrick (Ed.), *Contact metamorphism* (pp. 437–498). Chelsea, MI: Mineralogical Society of America.
- Gazel, E., Plank, T., Forsyth, D. W., Bendersky, C., Lee, C.-T., & Hauri, E. H. (2012). Lithosphere versus asthenosphere mantle sources at the Big Pine Volcanic Field, California. *Geochemistry, Geophysics, Geosystems*, 13, Q0AK06. <https://doi.org/10.1029/2012GC004060>
- Gonnermann, H., & Taisne, B. (2015). Magma transport in dikes. In H. Sigurdsson (Ed.), *The Encyclopedia of Volcanoes* (pp. 215–224). Cambridge: Academic Press.
- Gualda, G. A. R., Ghiorso, M. S., Lemons, R. V., & Carley, T. L. (2012). Rhyolite-MELTS: A modified calibration of MELTS optimized for silica-rich, fluid-bearing magmatic systems. *Journal of Petrology*, 53(5), 875–890. <https://doi.org/10.1093/petrology/egr080>
- Halliday, A. N., Fallick, A. E., Hutchinson, J., & Hildreth, W. (1984). A Nd, Sr, and O isotopic investigation into the causes of chemical and isotopic zonation in the Bishop Tuff, California. *Earth and Planetary Science Letters*, 68(3), 379–391. [https://doi.org/10.1016/0012-821X\(84\)90123-7](https://doi.org/10.1016/0012-821X(84)90123-7)
- Hetland, E. A., Simons, M., & Dunham, E. M. (2010). Postseismic and interseismic deformation due to fault creep I: Model description. *Geophysical Journal International*, 181(1), 81–98. <https://doi.org/10.1111/j.1365-246X.2010.04522.x>
- Hildreth, W. (2004). Volcanological perspectives on Long Valley, Mammoth Mountain, and Mono Craters: Several contiguous but discrete systems. *Journal of Volcanology and Geothermal Research*, 136, 169–198. <https://doi.org/10.1016/j.jvolgeores.2004.05.019>
- Hildreth, W., Fierstein, J., & Calvert, A. (2017). Early postcaldera rhyolite and structural resurgence at Long Valley Caldera, California. *Journal of Volcanology and Geophysical Research*, 335, 1–34. <https://doi.org/10.1016/j.jvolgeores.2017.01.005>
- Hildreth, W., & Wilson, C. J. N. (2007). Compositional zoning of the Bishop Tuff. *Journal of Petrology*, 48(5), 951–999. <https://doi.org/10.1093/petrology/egm007>
- Hofmeister, A. M. (1999). Mantle values of thermal conductivity and the geotherm from phonon lifetimes. *Science*, 283(5408), 1699–1706. <https://doi.org/10.1126/science.283.5408.1699>
- Holbrook, W. S., Mooney, W. D., & Christensen, N. I. (1992). The seismic velocity structure of the deep continental crust. In D. M. Fountain, R. Arculus, & R. W. Kay (Eds.), *Continental Lower Crust* (pp. 1–43). Amsterdam: Elsevier.
- Huang, W., & Russell, R. D. (2011). Adaptive moving mesh methods. *Applied Mathematical Sciences* (pp. 434). New York, NY: Springer. <https://doi.org/10.1007/978-1-4419-7916-2>
- Huppert, H. E., & Sparks, R. S. J. (1988). The generation of granitic magmas by intrusion of basalt into continental crust. *Journal of Petrology*, 29(3), 599–624. <https://doi.org/10.1093/petrology/29.3.599>
- Jackson, M. D., Cheadle, M. J., & Atherton, M. P. (2003). Quantitative modeling of granitic melt generation and segregation in the continental crust. *Journal of Geophysical Research*, 108(B7), 2332. <https://doi.org/10.1029/2001JB001050>
- Karakas, O., & Dufek, J. (2015). Melt evolution and residence in extending crust: Thermal modeling of the crust and crustal magmas. *Earth and Planetary Science Letters*, 425, 131–144. <https://doi.org/10.1016/j.epsl.2015.06.001>
- Kay, R. W., Kay, S. M., & Arculus, R. J. (1992). Magma genesis and crustal processing. In D. M. Fountain, R. Arculus, & R. W. Kay (Eds.), *Continental Lower Crust* (pp. 423–445). Amsterdam: Elsevier.
- Lange, R. A. (1997). A revised model for the density and thermal expansivity of K₂O-Na₂O-CaOMgO-Al₂O₃-SiO₂ liquids from 700–1900 K: Extension to crustal magmatic temperatures. *Contributions to Mineralogy and Petrology*, 130, 1–11. <https://doi.org/10.1007/s004100050345>
- Lange, R. A. (2013). The origin of highly evolved, voluminous rhyolites by progressive, multiple episodes of partial melting: The resolution of several paradoxes. *Geochemical Society Ingerson Lecture, Geological Society of America Abstracts with Programs*, vol. 45, no. 7, p. 84.
- Lange, R. A. (2019). Evidence for rapid transport of rhyolite melt during hydrous partial melting of granitoid: The critical role of pre-existing aplite dikes and fluid-undersaturated conditions. *American Geophysical Union Fall Meeting Abstracts*, V44A-01.
- Leeman, W. P., Annen, C., & Dufek, J. (2008). Snake River Plain—Yellowstone silicic volcanism: Implications for magma genesis and magma fluxes. *Geological Society, London, Special Publications*, 304, 235–259. <https://doi.org/10.1144/SP304.12>
- Levander, A., & Miller, M. S. (2012). Evolutionary aspects of lithosphere discontinuity structure in the western U.S. *Geochemistry, Geophysics, Geosystems*, 13, Q0AK07. <https://doi.org/10.1029/2012GC004056>
- Lim, S. H., Hetland, E. A., & Lange, R. A. (2012). Numerical models of transient partial melting of the lower crust during repeated emplacement of basalt sills and subsequent cooling due to advection of melt out of the lower crust. Abstract T31G-2708. Presented at 2012 Fall Meeting, AGU, San Francisco, Calif., 3–7 Dec.
- Manley, C. R., Glazner, A. F., & Farmer, G. L. (2000). Timing of volcanism in the Sierra Nevada of California: Evidence for Pliocene delamination of the basaltic root? *Geology*, 28(9), 811–814. [https://doi.org/10.1130/0091-7613\(2000\)28<811:TOVITS>2.0.CO;2](https://doi.org/10.1130/0091-7613(2000)28<811:TOVITS>2.0.CO;2)
- Metz, J. M., & Mahood, G. A. (1991). Development of the Long Valley, California, magma chamber recorded in precaldern rhyolite lavas of Glass Mountain. *Contributions to Mineralogy and Petrology*, 106(3), 379–397. <https://doi.org/10.1007/BF00324565>
- Noda, H., Dunham, E., & Rice, J. (2009). Earthquake ruptures with thermal weakening and the operation of major faults at low overall stress levels. *Journal of Geophysical Research*, 114, B07302. <https://doi.org/10.1029/2008JB006143>
- Ochs, F. A., & Lange, R. A. (1999). The density of hydrous magmatic liquids. *Science*, 283(5406), 1314–1317. <https://doi.org/10.1126/science.283.5406.1314>
- Oliver, H. W., Moore, J. G., & Sikora, R. F. (1993). Internal structure of the Sierra Nevada batholith based on specific gravity and gravity measurements. *Geophysical Research Letters*, 20(20), 2179–2182. <https://doi.org/10.1029/93GL01379>
- Petford, N., & Gallagher, K. (2001). Partial melting of mafic (amphibolitic) lower crust by periodic influx of basaltic magma. *Earth and Planetary Science Letters*, 193(3–4), 483–499. [https://doi.org/10.1016/S0012-821X\(01\)00481-2](https://doi.org/10.1016/S0012-821X(01)00481-2)
- Rudnick, R. L., & Gao, S. (2014). Composition of the continental crust. In *Treatise on Geochemistry* (2nd ed., Vol. 4, pp. 1–51). Amsterdam: Elsevier. <https://doi.org/10.1016/B978-0-08-095975-7.00301-6>
- Rushmer, T. (2001). Volume change during partial melting reactions: Implications for melt extraction, melt geochemistry and crustal rheology. *Tectonophysics*, 342, 389–405. [https://doi.org/10.1016/S0040-1951\(01\)00172-X](https://doi.org/10.1016/S0040-1951(01)00172-X)
- Saleeby, J., Ducea, M., & Clemens-Knott, D. (2003). Production and loss of high-density batholithic root, southern Sierra Nevada, California. *Tectonics*, 22(6), 1064. <https://doi.org/10.1029/2002TC001374>
- Shaw, H. R. (1980). Chapter 6. The fracture mechanisms of magma transport from the mantle to the surface. In *Physics of magmatic processes* (pp. 201–264). <https://doi.org/10.1515/9781400854493.201>

- Simakin, A. G., & Bindeman, I. N. (2012). Remelting in caldera and rift environments and the genesis of hot, "recycled" rhyolites. *Earth and Planetary Science Letters*, 337-338, 224-235. <https://doi.org/10.1016/j.epsl.2012.04.011>
- Simon, J. I., Weis, D., DePaolo, D. J., Renne, P. R., Mundil, R., & Schmitt, A. K. (2014). Assimilation of preexisting Pleistocene intrusions at Long Valley by periodic magma recharge accelerates rhyolite generation: Rethinking the remelting model. *Contributions to Mineralogy and Petrology*, 167(1), 1-34. <https://doi.org/10.1007/s00410-013-0955-5>
- Sleep, N. H. (1988). Tapping of melt by veins and dikes. *Journal of Geophysical Research*, 93(B9), 10,255-10,272. <https://doi.org/10.1029/jb093ib09p10255>
- Spera, F. J., & Bohrsen, W. A. (2018). Rejuvenation of crustal magma mush: A tale of multiply nested processes and timescales. *American Journal of Science*, 318, 90-140. <https://doi.org/10.2475/01.2018.05>
- Stachnik, J. C., Dueker, K., Schutt, D. L., & Yuan, H. (2008). Imaging Yellowstone plume-lithosphere interactions from inversion of ballistic and diffusive Rayleigh wave dispersion and crustal thickness data. *Geochemistry, Geophysics, Geosystems*, 9, Q06004. <https://doi.org/10.1029/2008GC001992>
- Thurber, C., Zhang, H., Brocher, T., & Langenheim, V. (2009). Regional three-dimensional seismic velocity model of the crust and uppermost mantle of northern California. *Journal of Geophysical Research*, 114, B01304. <https://doi.org/10.1029/2008JB005766>
- Wallace, P. J., Anderson, A. T., & Davis, A. M. (1999). Gradients in H₂O, CO₂, and exsolved gas in a large-volume silicic magma system: Interpreting the record preserved in melt inclusions from the Bishop Tuff. *Journal of Geophysical Research*, 104(B9), 20,097-20,122. <https://doi.org/10.1029/1999JB900207>
- Waters, L. E., & Lange, R. A. (2017). Why aplites freeze and rhyolites erupt: Controls on the accumulation and eruption of high-SiO₂ (eutectic) melts. *Geology*, 45(11), 1019-1022. <https://doi.org/10.1130/g39373.1>
- Wilson, C. J. N., & Hildreth, W. (1997). The Bishop Tuff: New insights from eruptive stratigraphy. *The Journal of Geology*, 105(4), 407-440. <https://doi.org/10.1086/515937>
- Yuan, H., Dueker, K., & Stachnik, J. (2010). Crustal structure and thickness along the Yellowstone hot spot track: Evidence for lower crustal outflow from beneath the eastern Snake River Plain. *Geochemistry, Geophysics, Geosystems*, 11, Q03009. <https://doi.org/10.1029/2009GC002787>
- Zandt, G., Hersh, G., Owens, T. J., Ducea, M., Saleeby, J., & Jones, C. H. (2004). Active foundering of a continental arc root beneath the southern Sierra Nevada in California. *Nature*, 431(7004), 41-46. <https://doi.org/10.1038/nature02847>
- Zhang, Y., Xu, Z., Zhu, M., & Wang, H. (2007). Silicate melt properties and volcanic eruptions. *Reviews of Geophysics*, 45, RG4004. <https://doi.org/10.1029/2006RG000216>